

Chasing the Unicorn: RHIC and the QGP

Unicorn = fantastic and mythical beast!

RHIC = Relativistic Heavy Ion Collider @ Brookhaven Natl. Lab (BNL):
collide large nuclei at high energies (also: SPS & LHC @ CERN)

QGP = Quark Gluon Plasma =
New state of hadronic matter, in
thermodynamic equilibrium at temperature $T \neq 0$

Q: Has RHIC made the QGP?

1. QCD @ nonzero temp.: the QGP
2. The QGP on the Lattice: numerical “experiment”
3. Experiments at RHIC: evidence for “gluon stuff” -
the (high-pt) tail wags the (low-pt) body of the Unicorn



A: Some new kind of matter has been created

Quark Model: Particle Zoo. Masses, Scales....

Hadronic particles = strongly interacting = *baryons & mesons*

baryons: proton & neutron: mass = 940 MeV (MeV = 10^6 eV) \approx 1 GeV

mesons: pions, π^\pm, π^0 mass $\pi \approx$ 140 MeV

All hadrons interact by pion exchange \Rightarrow fund. length: $1/m_\pi =$ 1 fermi (fm)
fund. time scale = 1 fm/c

Less familiar: **strange baryons:** Λ (1120), Σ (1190), Ω (1680)

strange mesons: 4 Kaons (K^\pm, K^0, \bar{K}^0) (540) & η (550) (η' (980)?)

Above: mesons spin 0, baryon's spin 1/2.

Also: spin 1 mesons: not strange, ρ (770), and strange ϕ (1020).

Ignore heavier particles, such as J/Ψ , Υ ...

Quark Model: all hadrons composed of quarks.

Quark Zoo: 2, 3, 5(!) quarks

Above hadrons: from **up, down, & strange quarks** = u, d & s: **3 quark flavors**
(Heavier flavors: charm, bottom, & top quarks)

π , K, η very light = (approx.) “spin waves” of (approx.) **chiral symmetry**
mass u&d (5,10 MeV) \ll mass s quark (100 MeV)

mesons = $\bar{q}q$ $q = u, d, \text{ or } s \text{ quark.}$ $\pi = u\&d$ $K = (u \text{ or } d) \& s$

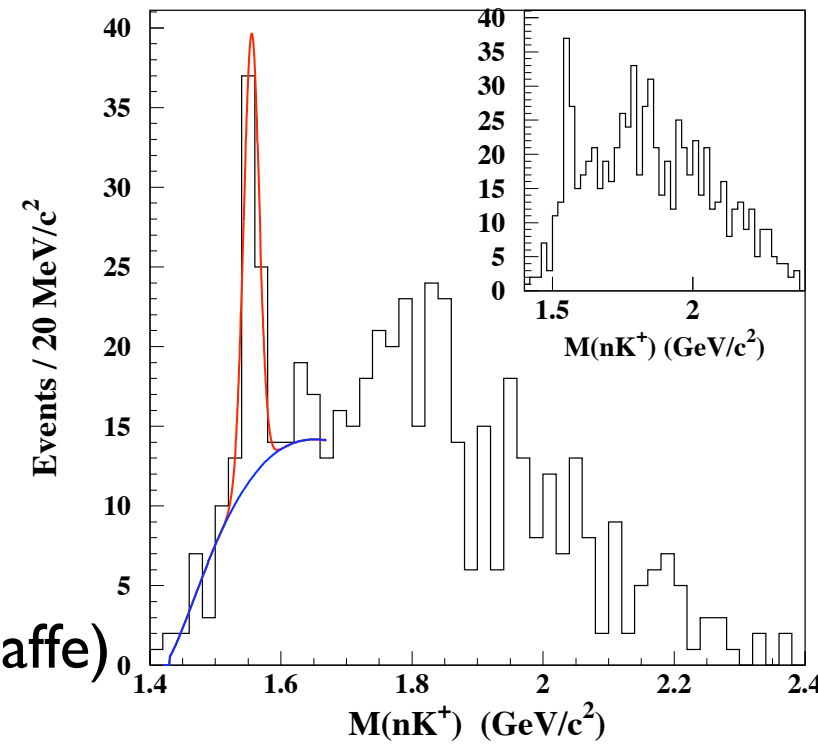
baryons = qqq $N = u,d's$; $\Lambda = 2(u,d)\&s$; $\Sigma = (u,d)\& ss$; $\Omega = sss$.

New: “penta-quarks” = $qqqq\bar{q}$

$\Theta^{++} = uud\bar{s}$ (Diakonov, Petrov, Polyakov)

CLAS @ JLAB: mass = $1150 \pm 10 \text{ MeV}$. **8σ !**
width < 26 MeV! \Rightarrow really narrow

Where is the (hexa-quark) H-dibaryon = uuddds? (Jaffe)



QCD: Quark Model + *Gluons*

Global symmetries familiar. E.g., spherical symmetry = $SO(3)$.
Uniform rotation everywhere the same in space.

Local symmetry: at *each* point, *independent* rotations in “internal” space.

Need new degrees of freedom: non-Abelian gauge fields = *gluons*.

QCD = $SU(3)$ gluons + quarks. 3 of $SU(3)$ = # colors.

Analogy: for critical point in *four* dimensions (= critical dim.), $\lambda \phi^4$ int.

$$\lambda(p) \sim \frac{1}{\log(\mu/p)} \quad p \Rightarrow 0 \quad \begin{array}{l} p = \text{momentum, } \mu = \text{ren. mass scale} \\ \Rightarrow \text{compute in } \lambda \text{ as } p \Rightarrow 0, \text{ large distances} \end{array}$$

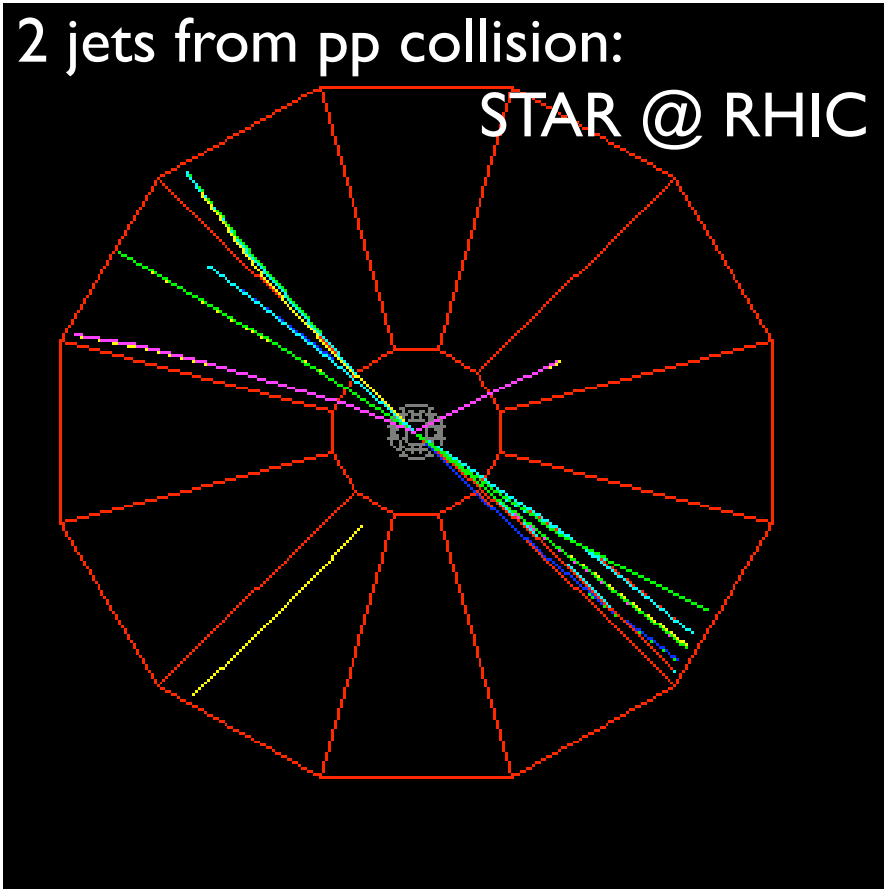
QCD = converse = *asymptotically free*: g^2 = QCD coupling constant

$$g^2(p) \sim \frac{1}{\log(p/\mu)} \quad p \Rightarrow \infty \quad \begin{array}{l} \Rightarrow \text{compute in } g^2 \text{ at large momentum} \\ = \text{short distances} \end{array}$$

Jets: “seeing” quarks and gluons in QCD

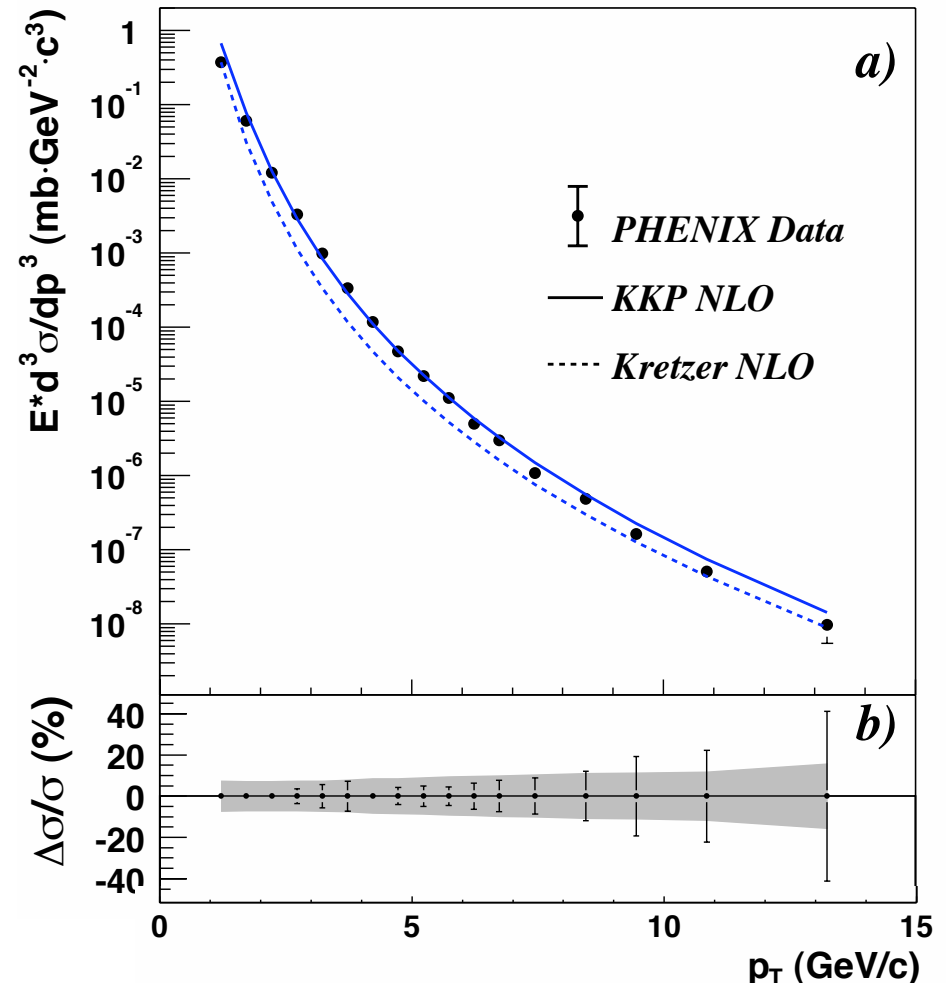
2 jets from pp collision:

STAR @ RHIC



Jets can be reliably computed in perturbation theory, down to momenta \sim few (5, 50?) GeV!
At high energies, can tell, *indirectly*, between gluon jets & quark jets:
on average, gluon jets are “fatter”.

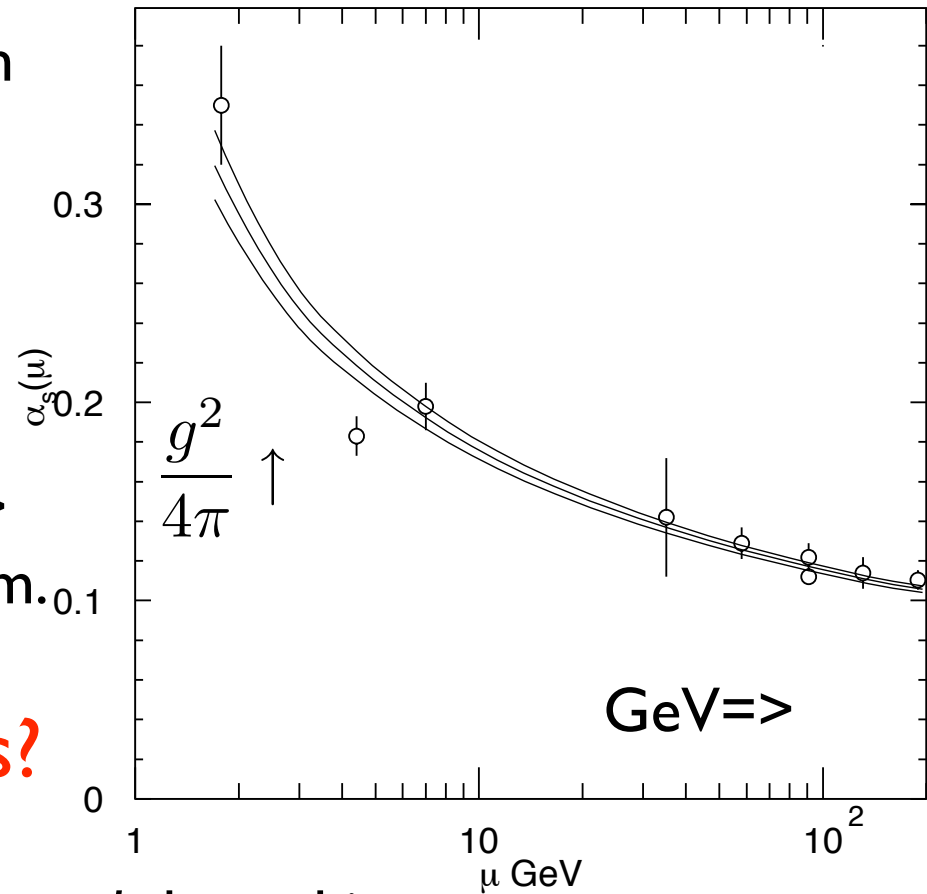
At high energies, energetic quarks and gluons produce **jets**. While rare, they are a striking feature. Note: by momentum conservation, any jet in one direction has a *backward* jet in the opposite direction.



QCD @ Low Energies: Confinement

For a critical point at the critical dimension the coupling vanishes at small momentum, grows at large mom. ("Landau's ghost")

Conversely, in QCD, while the coupling is small at high momentum, it grows as the momentum decreases. To the right: => variation of QCD coupling with momentum.



What happens at large distances?

Confinement: quarks and gluons are *permanently* bound into **color-neutral** states = mesons and baryons.

“Infrared slavery”: **linear potential between quarks** (more later)

Q: where are the states with pure glue = **glueballs**?

How to see the “glue” at large distances?

Symmetries of QCD: Chiral Symmetry

Most “familiar”: *chiral symmetry*.

Like a magnet: *broken* at low temperature, restored at some finite temperature.

In broken phase, (approx.) “spin waves” = (almost massless) pions

up & down quarks: “*flavor*” symmetry = $SU_L(2) \times SU_R(2) = O(4)$

L(ef), R(ight) = *chirality*, *special* to massless fermions.

$O(4)$ vector = $(\sigma, \vec{\pi})$. At zero temp, condensate: $\langle \sigma \rangle \sim \langle \bar{q}q \rangle \neq 0$

With strange quark, *flavor symmetry* = $SU_L(3) \times SU_R(3)$

$\langle \sigma \rangle \sim \langle \bar{q}q \rangle \neq 0 \Rightarrow 3 \pi$'s, 4 K 's, 1 η are massless. Correct # Goldstone bsns

(What about η' from extra axial $U(1)$? Instantons....

Could dramatically affect transition properties with *light* quarks.)

Deconfinement as a *Global* $Z(3)$ Symmetry

't Hooft: **rigorous** order parameter for **confinement** .

Consider multiplying each quark by a **constant** phase:

$$q \rightarrow e^{2\pi i/3} q \quad , \quad \bar{q} \rightarrow e^{-2\pi i/3} \bar{q}$$

Mesons and baryons are invariant under this **global transformation**:

$$\bar{q}q \rightarrow \bar{q}q \quad , \quad qqq \rightarrow (e^{2\pi i/3})^3 qqq = qqq$$

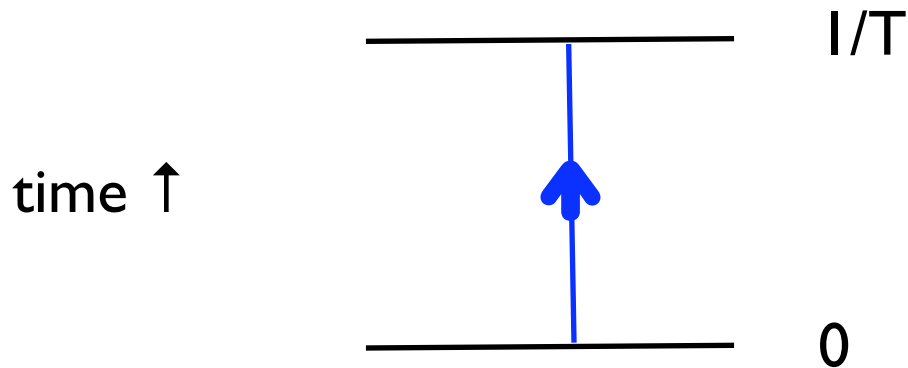
but any other states, such as q , qq , etc, are not. We could also use $e^{-2\pi i/3}$ as well, but only these transformations (and I !) are allowed. This is a global symmetry of $Z(3)$ (the third roots of unity).

Hence: **confinement = unbroken global symmetry of $Z(3)$** .

But! Only valid in a **pure** gauge theory, **without** dynamical quarks. These quarks above are “**test**” quarks. Dynamical quarks act as sources of $Z(3)$ flux, and spoil the symmetry. In QCD, is the $Z(3)$ symmetry **approximate**?

Test Quarks & Polyakov Loops

How to construct a test quark? Consider a nonzero temperature T : in imaginary time formalism, euclidean time runs from 0 to $1/T$. Put an *infinitely* heavy quark down at some point in space: all it can do is run up in time:



While this test quark can't move, it *can* exchange *color* with the thermal bath. It does this through a (color) Aharonov-Bohm phase factor:

$$\ell = \frac{1}{3} \text{tr} \mathcal{P} \exp \left(ig \int_0^{1/T} A_0 d\tau \right)$$

= (trace of) propagator for test quark. Known as **Polyakov loop**, or trace of (thermal) **Wilson line**. (Wraps around in imaginary time \Rightarrow loop).

Deconfinement & Polyakov Loops

't Hooft: part of *local* SU(3) is *global* Z(3) (not obvious!) $\ell \rightarrow e^{2\pi i/3} \ell$

At T=0, *confinement* => quarks don't propagate => *UNbroken* Z(3) symmetry

$$\langle \ell \rangle = 0 \quad , \quad T < T_{deconf}$$

As $T \rightarrow \infty$, by *asymptotic freedom*, coupling g^2 is small, pert. thy. ok. So the global Z(3) symmetry is (spontaneously) *broken*:

$$\langle \ell \rangle \neq 0 \quad , \quad T > T_{deconf}$$

Hence there is a temperature at which the loop gets a v.e.v.:

$T_d \equiv T_{deconf}$ = temperature for the *deconfining* phase transition.

Deconf. opposite to spins: Z(3) broken at high, and not low, temp.

In terms of order of the transition, just like typical Z(3) spins.

Order of Phase Transitions

Most cases follow from simple mean field analysis:

Deconfining transition: cubic invariant is $Z(3)$ symmetric: ℓ^3
 \Rightarrow *first order deconfining trans.* (Svetitsky & Yaffe).

Chiral transition: for **two** massless flavors, $O(4)$ sym. \Rightarrow *second order chiral trans*
For **three** massless flavors, $SU_L(3) \times SU_R(3)$ symmetry. Again, cubic invariant $\det(\Phi) \Rightarrow$ *first order chiral transition*.

Tech.'y: no restoration of axial $U(1)$. Even if so, still *first order chiral transition*:
“fluctuation-induced” first order, like superconductor (RDP & Wilczek).

Guess: First Order Transition(s)?

“Of course”! Hadrons \neq Quarks & Gluons.

But: relation between deconfining and chiral transitions? 1 or 2 transitions?
For QCD, both $Z(3)$ and chiral symmetries are *approximate*.

The “Unicorn”:

Quark-Gluon Plasma =

Deconfined,
Chirally Symmetric “Phase”

But how to compute
properties of the QGP?



QGP on the Lattice: “Numerical Experiment”

How to compute properties of transition = *strong coupling regime*?

Put theory on lattice! Then simulate *numerically*.

Wegner & Wilson: *easy* for *local* gauge theory: *quarks on sites, gluons on links*.

$$\begin{array}{c} \bar{q} \left| \begin{array}{c} \hat{n}^\mu \rightarrow \\ \hline U = e^{igaA_\mu \hat{n}^\mu} \end{array} \right| q \end{array}$$

Lattice spacing = a . *Asymptotic freedom* \Rightarrow *unique* result for $a \Rightarrow 0$ ($p \Rightarrow \infty$)

Example of *universality*: e.g., at 2nd order transition, over large distances, *critical exponents unique* (func. of symmetry group & dimension)

Here, “*dimensional transmutation*”:

once value of coupling is fixed at some scale, *nothing* left to fix.

All dimensionless ratios are unique (func. of symmetry group & dimension)

What the Lattice can do

But: how close is the lattice (today) to the continuum limit, $a=0$?

“Pure” gauge (no dynamical quarks): present methods close to $a=0$!

QCD: present methods *not* close to $a=0$. All results tentative.

Very hard to put *global* chiral symmetry on lattice!

View: lattice simulations as (another) experiment... What it has told us so far:

Pure gauge: pressure for *all* temp.'s. $T_d \sim 270 \pm 10$ MeV.

Weakly first order deconfining trans.

Non-perturbative QGP from $T_d \Rightarrow 3 T_d$. NO “Of Course”

With quarks: $T_c \sim 175 \pm ?$ MeV

Order? Crossover today.

Only *one* transition (chiral = deconfining)

“Flavor independence”: pressure *with* qks like that *without* qks.

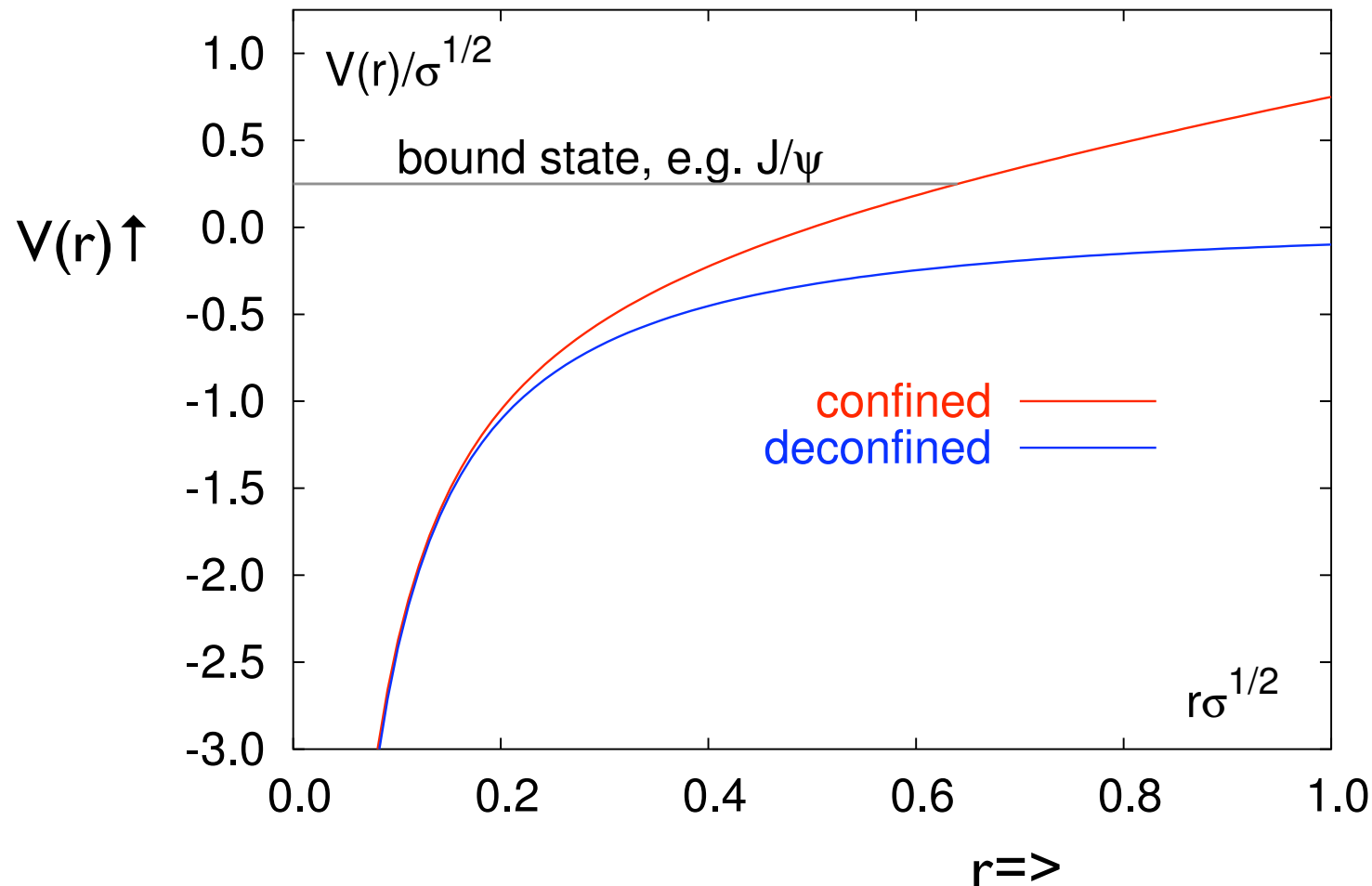
Confinement and the Quark Potential

$T < T_c$: $V(r)$ = potential between test quark and anti-quark,

$V(r) \sim \sigma r$ as $r \rightarrow \infty$: σ = string tension.

Linearly rising potential \Rightarrow permanent confinement

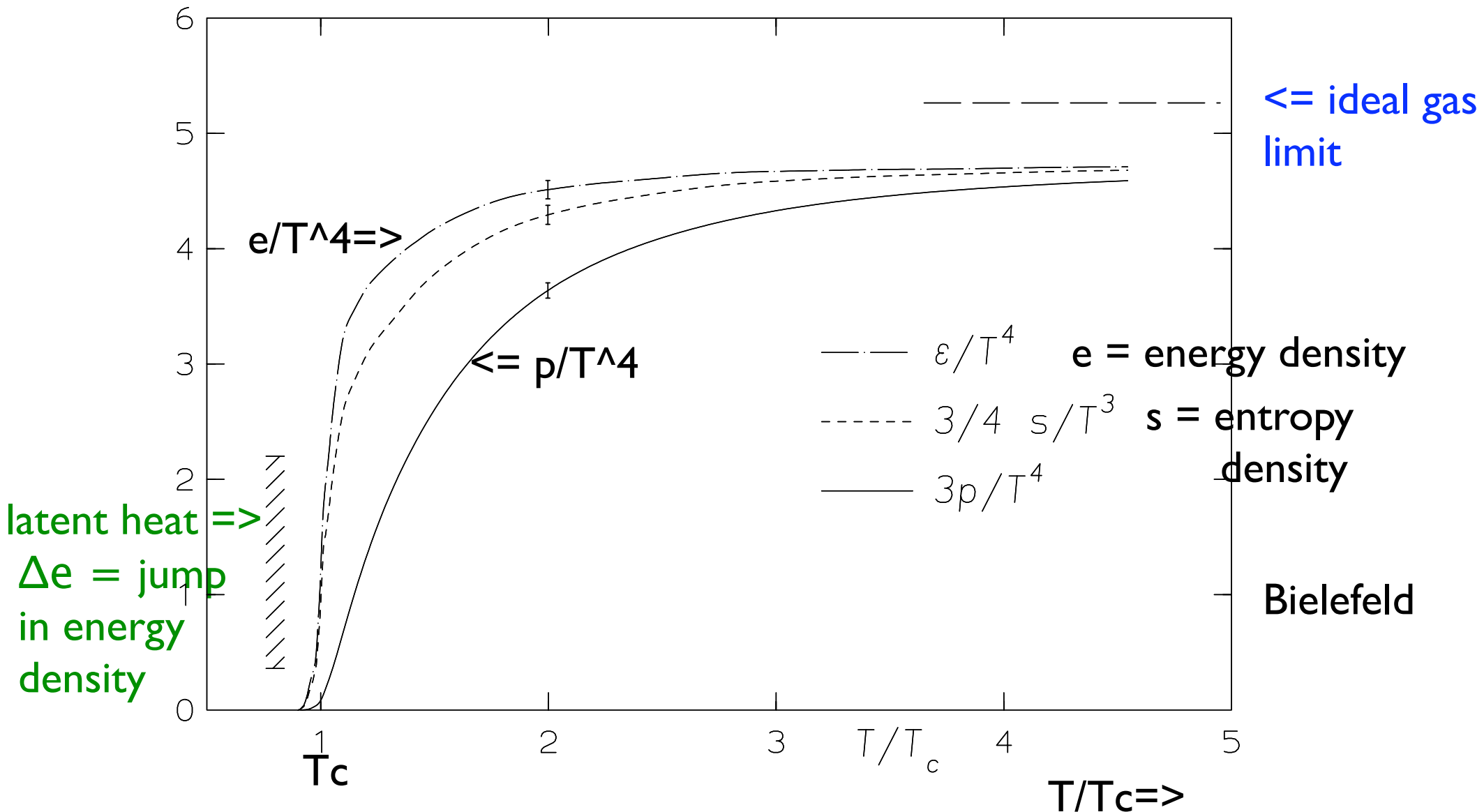
$T > T_c$: $V(r) \rightarrow \text{constant}$ as $r \rightarrow \infty$. Deconfinement.



Lattice: “Pure” Gauge Thermo. for SU(3)

In equilibrium, everything follows from pressure $p=p(T)$ (T =temperature)

Asymptotic freedom $\Rightarrow p/T^4 \Rightarrow$ constant as $T \rightarrow \infty$.



Lattice: “Pure” Gauge SU(3) Thermo.

Find: can extrapolate to continuum limit reliably.

$T_c \sim 270 \text{ MeV} \pm 5 \%$ (scale - and error! - dom.'d by string tension)

$T < T_c$, pressure very small in the confined phase.

Pure gauge \Rightarrow spectrum *massive* glueballs. Lightest glueball $\sim 1.5 \text{ GeV}$

Pressure of *heavy* glueballs $\sim \exp(-m/T_c) \ll 1$ very small. Or: T_c small.

$T \rightarrow \infty$: asymptotic freedom $\Rightarrow p \rightarrow$ ideal gas as $T \rightarrow \infty = 2 \times 8 \times \frac{\pi^2}{90} T^4$

$T > T_c$: relatively rapid approach to ideal gas: $\sim 80\%$ ideal gas by $3 T_c$.

Suggests: *non-perturbative* behavior: $T_c \Rightarrow 3 T_c$, “semi”-pert. $> 3 T_c$.

‘81 \Rightarrow ‘89: coarse lattices, far from continuum limit: strongly first order.

$>$ ‘89: deconfining transition *weakly* first order (APE, Columbia).

(Some) correlation lengths grow by ~ 10 near T_d !

SU(3) “Pure” Gauge: *Weakly* First Order

Latent heat: $\Delta e / e_{\text{ideal}} \sim 1/3$ (vs $4/3$ in bag model). So?

Look at physical correlation lengths, related to two point function of loops

$T < T_c$

string
tension

$$\frac{\sigma(T)}{\sigma(0)}$$

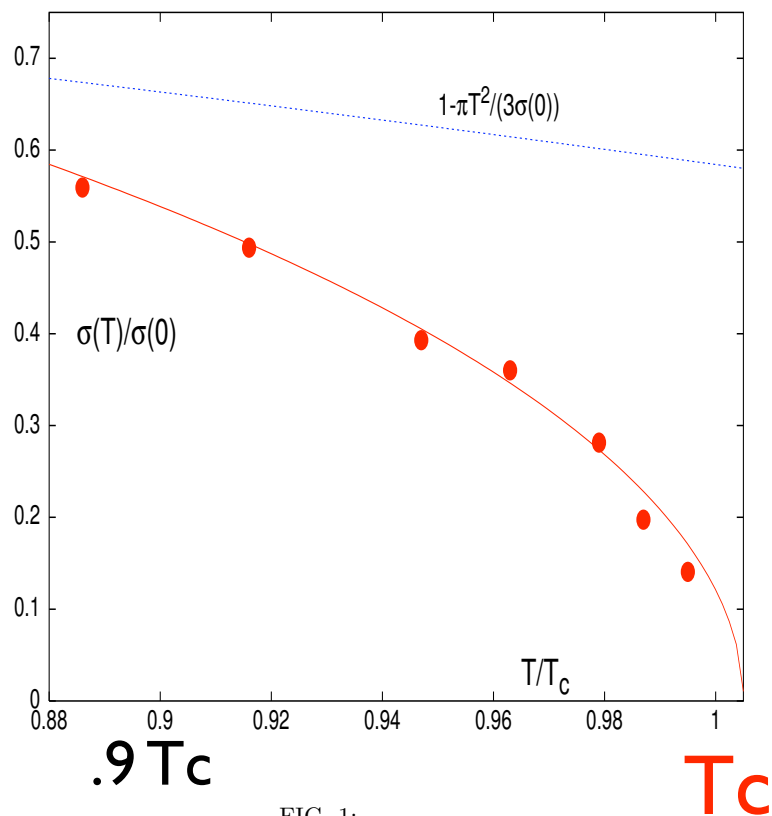


FIG. 1:

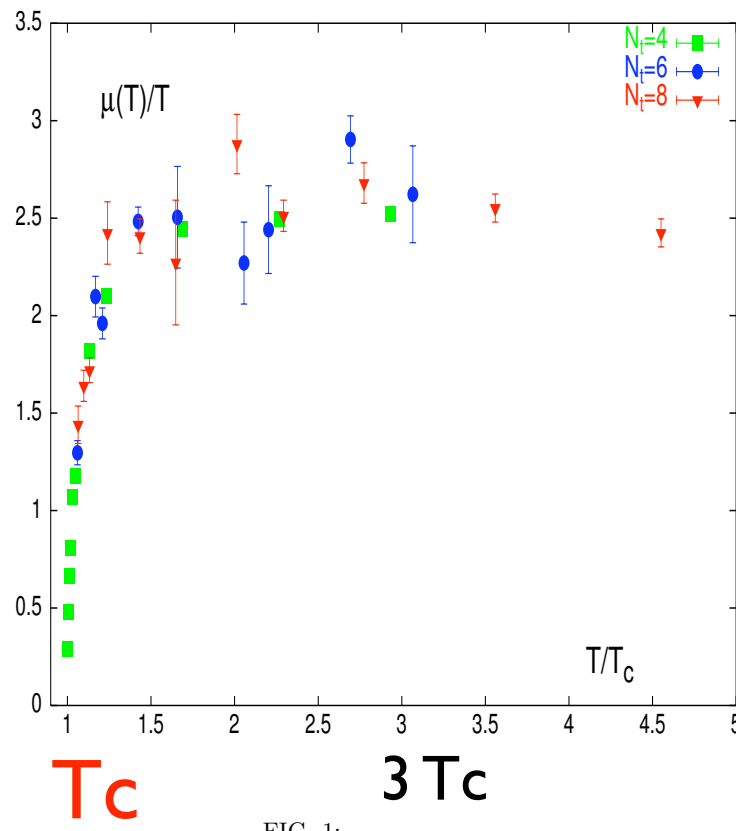


FIG. 1:

$T > T_c$

Debye
mass

$$\frac{m_{\text{Debye}}(T)}{m_{\text{Debye}}(0)}$$

Bielefeld

(Some) physical correlation
lengths grow by ~ 10 !

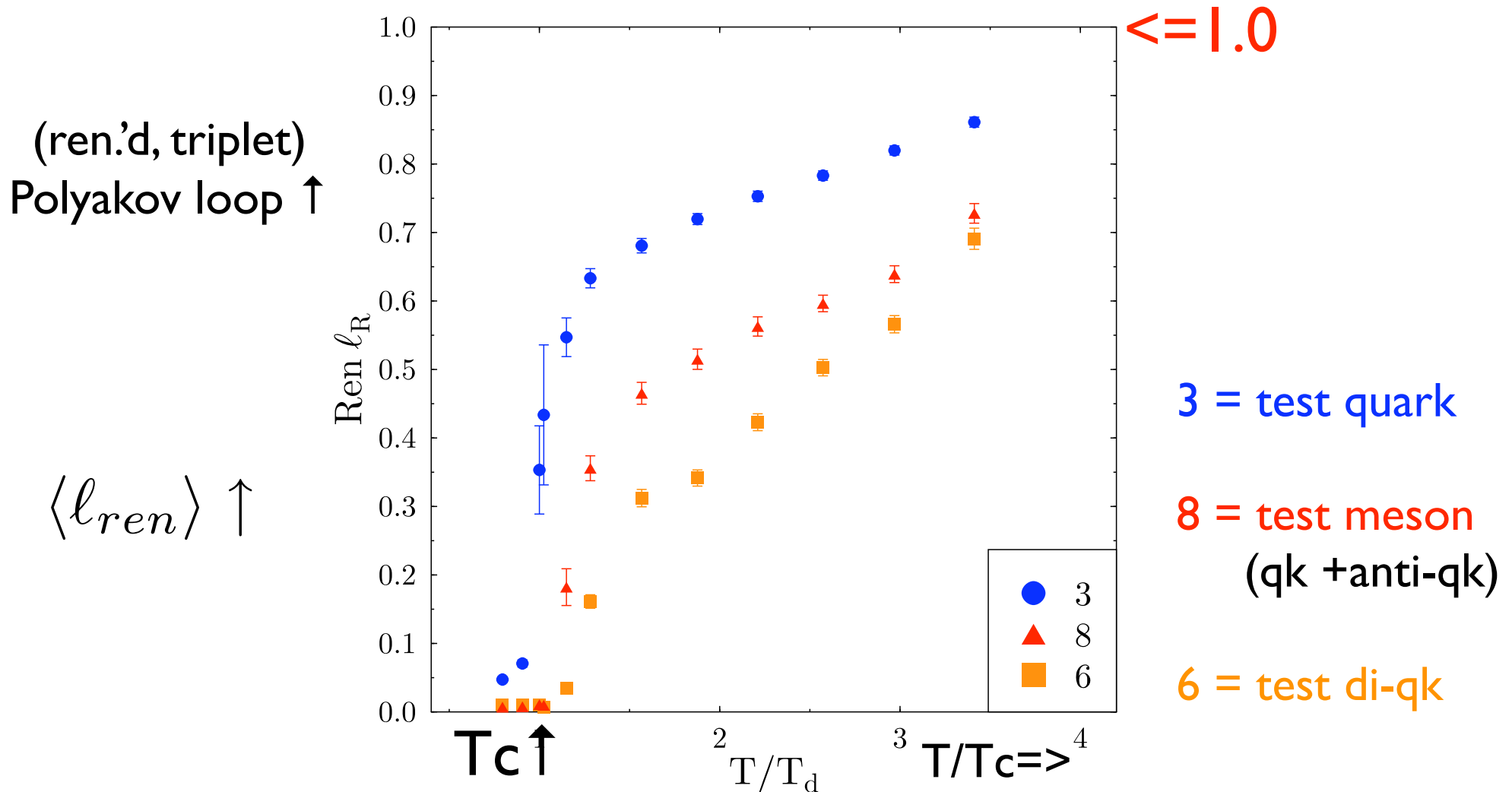
$$\frac{\sigma(T_d^-)}{\sigma(0)} \approx \frac{m_{\text{Debye}}(T_d^+)}{m_{\text{Debye}}(1.5 T_d)} \approx \frac{1}{10}$$

Lattice: *Non-pert.* QGP: $T_d \Rightarrow 3 T_c$

In pert. theory, ren'd (triplet) loop $\sim 1 + \dots$

(Ren'd) triplet loop < 1 for $T_c \Rightarrow 3 T_c \Rightarrow$ *non-perturbative* QGP $\leq 3 T_c$

Also: persistence of bound states...



Dumitru, Hatta, Lenaghan, Orginos, & RDP. Also: Bielefeld.

Lattice Thermo: *Big changes with Quarks*

QCD: “2+1” flavors (up & down light, strange heavy):

$T \rightarrow \infty$: ideal gas limit increases by ~ 3 .

pure glue: 16 times pressure for massless boson $= \pi^2 T^4 / 90$

3 massless flavors: 48.5 times pressure for 1 boson = ...

Still: pressure rapidly approaches ideal gas by ~ 3 times T_c

T_c : decreases by ~ 2 . Assured: T_c decreases as # flavors increases.

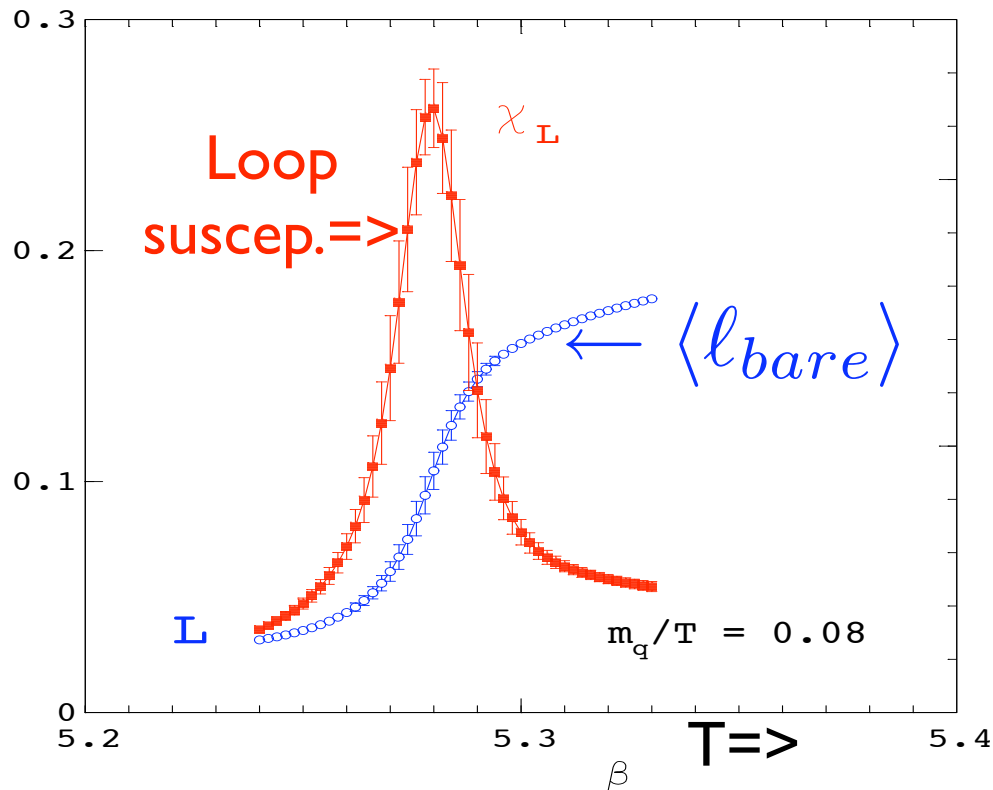
QCD: $T_c \sim 175 \pm ?$ MeV

Not close to continuum limit; hard getting quarks light enough (state of art: kaons ok, pions not)

$T < T_c$: in “confined” phase with pions (chiral symmetry broken),
pressure *small*: turns on only near T_c .

Lattice: Always ONE Phase Transition!

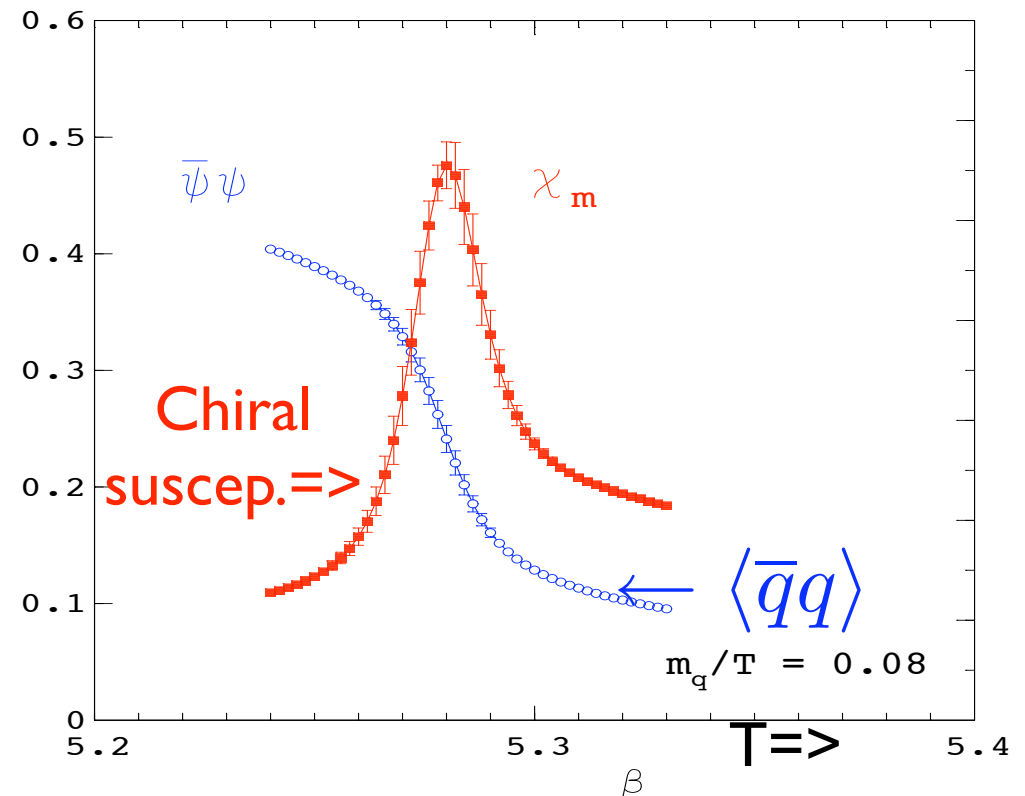
Could be two transitions, deconfining and chiral. NEVER seen for any quark mass.



=> chiral order parameter vs lattice coupling \sim temperature.
Also: chiral susceptibility.

Both susceptibilities peak at SAME temperature!

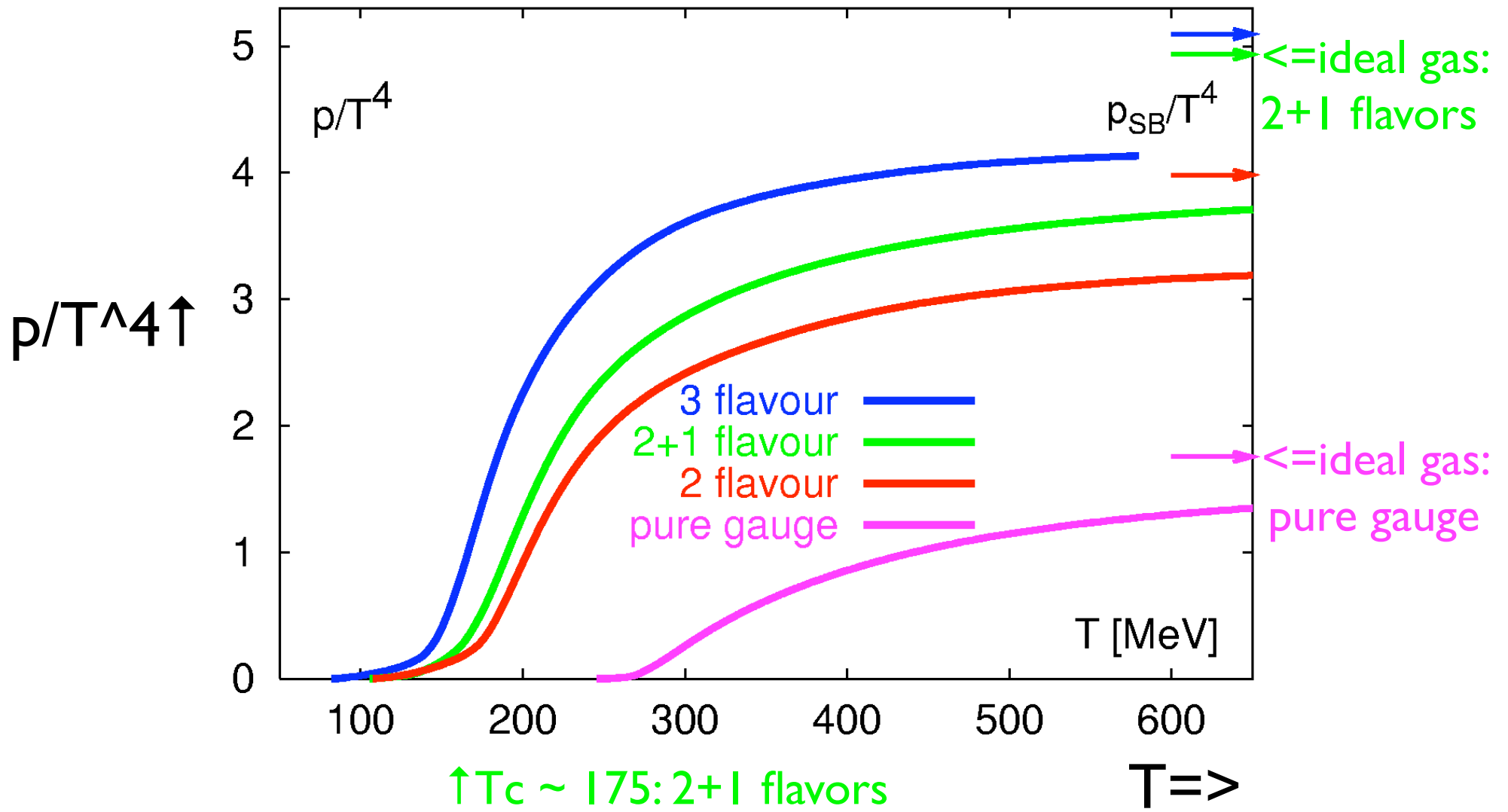
\leq (bare) Polyakov loop vs lattice coupling \sim temperature.
Also: loop susceptibility.



Lattice: Pressure vs T, Different # Flavors

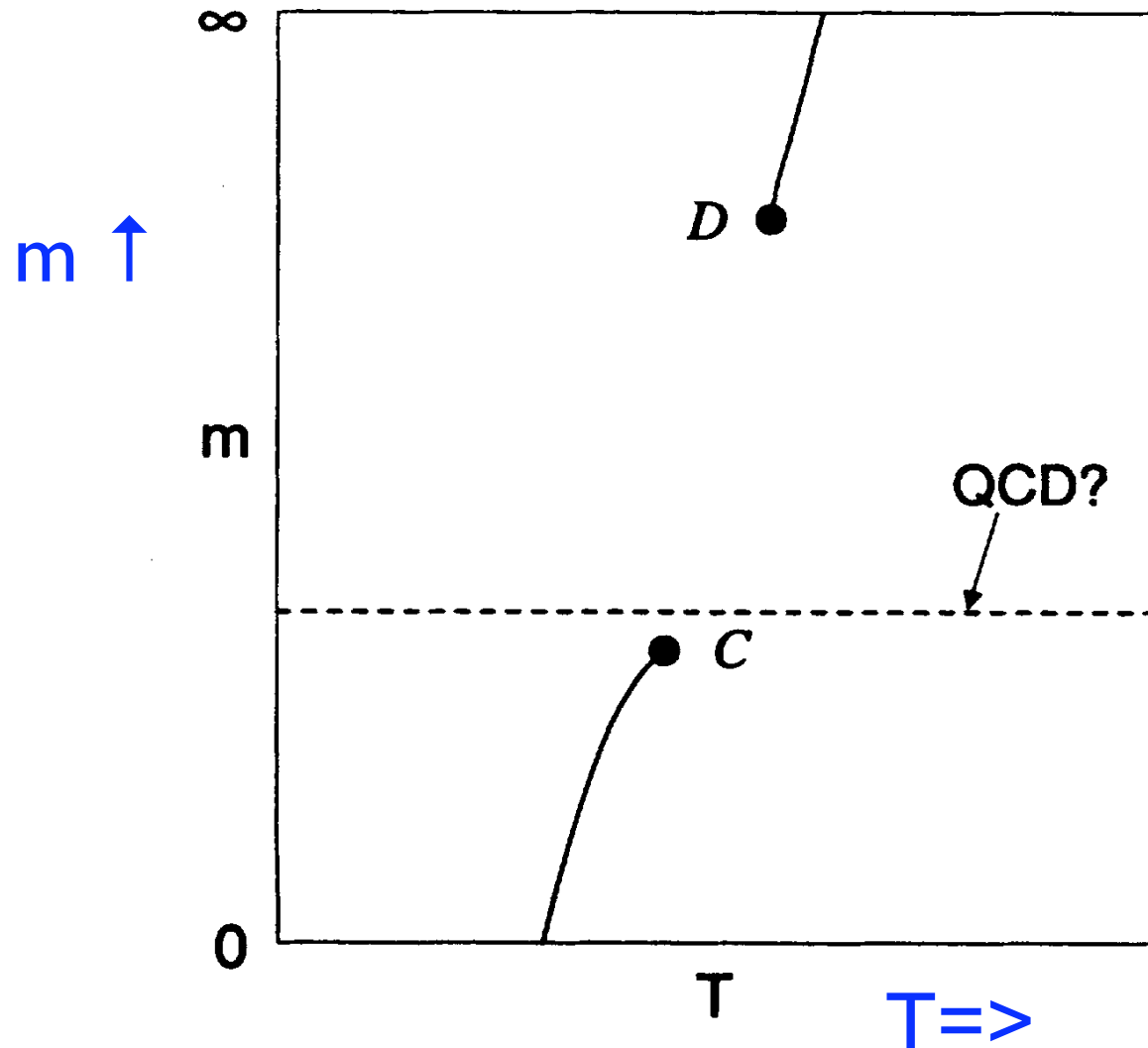
QCD: “2+1” flavors (up & down light, strange heavy): **BIG changes**

$p=p(T)=\text{pressure}$. Plot p/T^4 , \Rightarrow constant as $T \rightarrow \infty$ (asymptotic freedom)



Lattice: Order *VERY* Sensitive to Quarks

“Columbia” phase diagram: keep up, down, strange quarks in **fixed** ratio, vary overall mass scale = m .



First order for:

pure gauge ($m = \infty$)
3 massless flavors.

But deconfining (D) and
chiral (C) critical end-points!

Today:

QCD \sim *crossover*

\Rightarrow No phase transition.

True today. Role of axial $U(1)$?

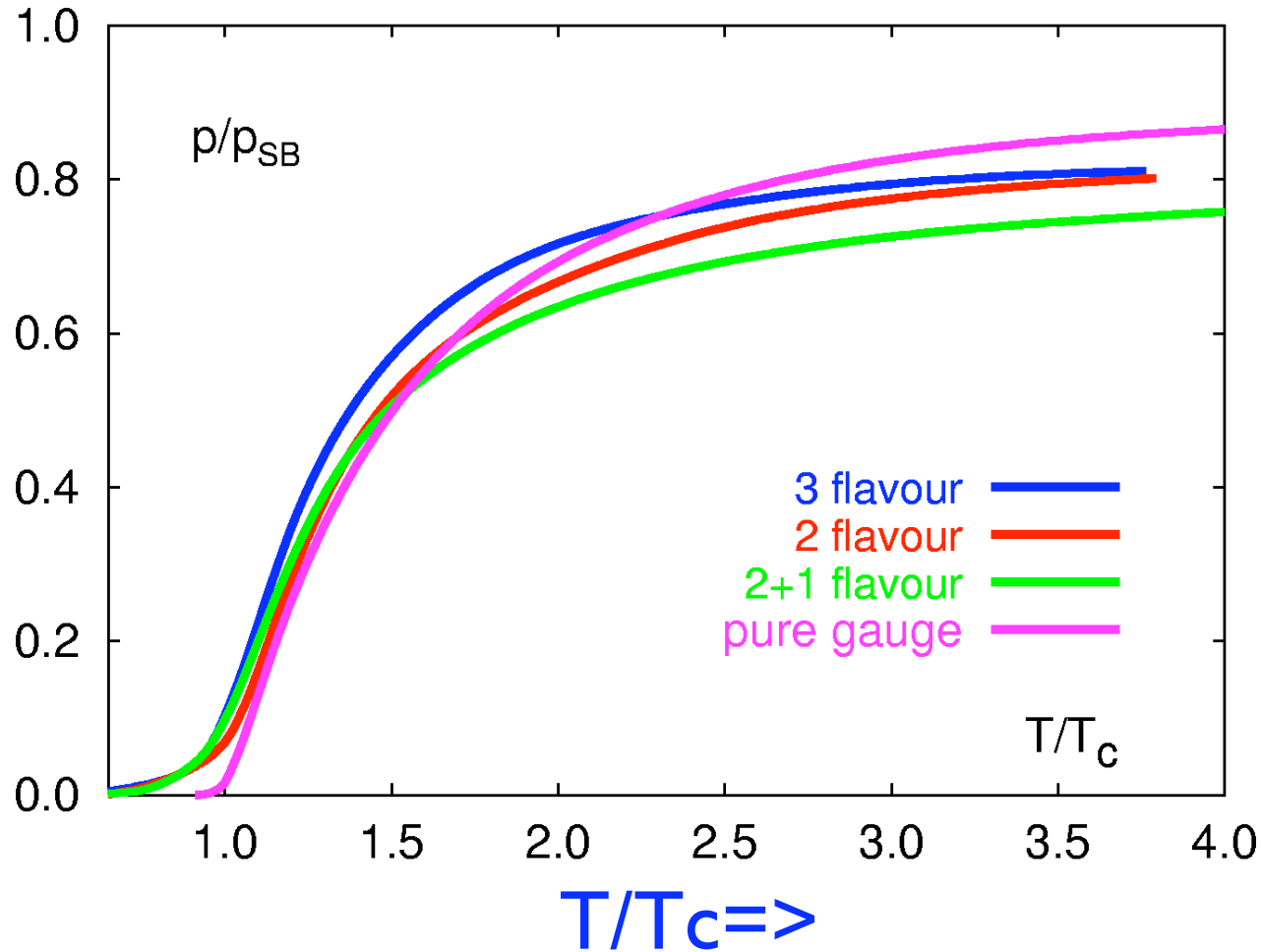
Lattice: “Flavor Independence”

Lattice finds *amazing* property:
properly scaled, pressure *with* quarks
like that *without*: *Bielefeld*.

$$\frac{p}{p_{ideal}} \left(\frac{T}{T_c} \right) \approx \text{universal}$$

1.0=>

pressure/
ideal gas ↑

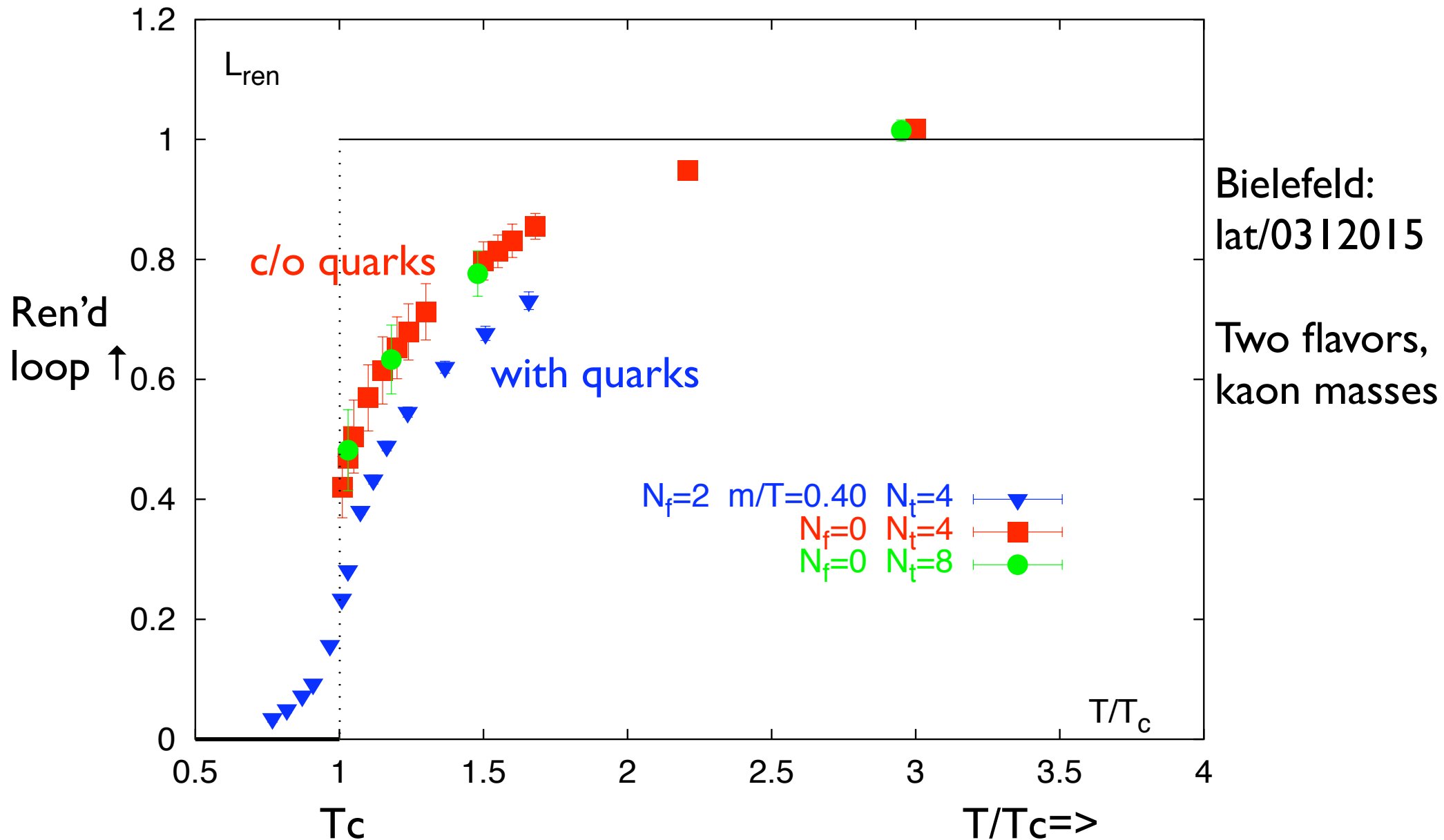


=> pressure
dominated by
gluons?

(Ren.'d) Polyakov Loop with Quarks ~ Pure Gauge

(Ren.'d) Loop approx.'y same with quarks as without =>

pressure dominated by gluons? (= Polyakov loop)



Hunting for the “Unicorn”: the Quark-Gluon Plasma, in Heavy Ion Collisions



“Unicorn” & the QGP: Scott, Stock, Gyulassy...

Why do AA? Big Transverse Size.

First, some essential definitions. One can collide:

pp: protons on protons. Serves as benchmark for “ordinary” hadronic coll’s.

AA: nucleus with **atomic number A** on the same type of nucleus.

pA: proton on a nucleus. At RHIC, often **dA** ($d = n+p \sim p$) for accelerator reasons (charge/mass ratio) **Serves as test to tell pp from AA.**

WHY AA? Nucleon’s are like hard spheres, so nuclear size $r_A \sim A^{1/3}$

Biggest: **Pb** (lead) or **Au** (gold), $A \sim 200 \Rightarrow r_A \sim 7$.

Transverse radius of nucleus $\sim A^{2/3} \Rightarrow$ **trans. size ~ 50 x proton.**

$A \rightarrow \infty$: infinite nuclear matter. $A \sim 200$ close to ∞ ? **Decide by experiment.**

(Very) roughly: transition from p to large A for $A \sim 30-50$.

Colliders: Energy, Machines

Particles accelerated in rings. Highest energy is for two rings, with particles travelling in opposite directions = *collider*.

Basic invariant: **total energy in the center of mass**, $E_{c.m.} \equiv \sqrt{s}$

For AA collisions, energy *per* nucleon is $\sqrt{s}/A \equiv \sqrt{s_{NN}}$

Machines:

$$\sqrt{s}/A$$

SPS @ CERN:	5 => 17 GeV	(fixed target)
RHIC @ BNL:	20, 130, 200 GeV	(collider)
LHC @ CERN:	5500 GeV = 5.5 TeV	(collider, > 2007)

SPS = Super Proton Synchrotron: CERN @ Geneva, Switzerland.

RHIC = Relativistic Heavy Ion Collider; BNL @ Long Island, NY

LHC = Large Hadron Collider.

Collider Kinematics

At low energies, form one “blob” which is radially symmetric = Landau model.

At high energies ($s \gg 1 \text{ GeV}$) particles go through each other. Use:

Momenta transverse to the beam: p_t

Momenta along the beam = p_z Exp.'y, not useful. Instead:

Rapidity = y : $y = \log\left(\frac{E + p_z}{E - p_z}\right)$ $y=0 = 90^\circ$ for collider

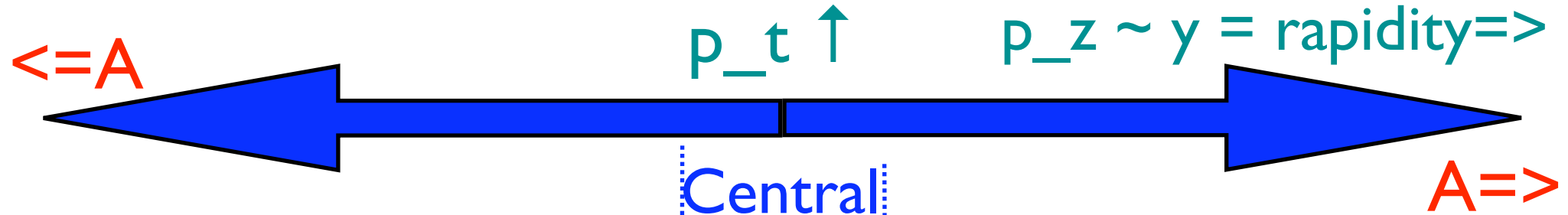
Pseudo-rapidity: η If one doesn't have particle ID, so assume $E = \sqrt{p_t^2 + p_z^2}$

Usually: # particles vs p_t , & y : *most particles at zero p_t , zero y .*

“Central regime:” “free” of incident nucleons, rapidity $y \sim 0$
=> most likely to exhibit $T \neq 0 \approx$ small net baryon density

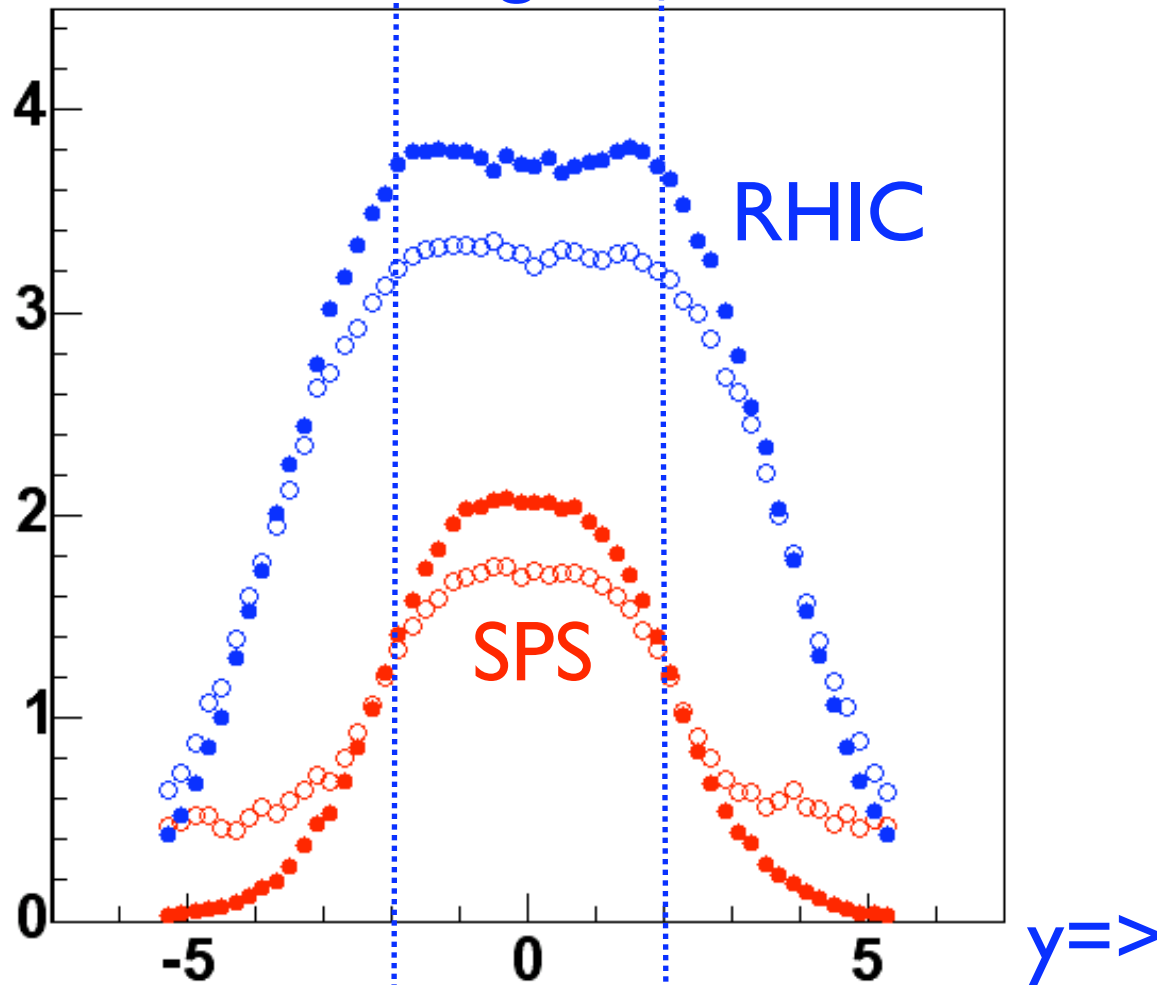
“Fragmentation regime”: where incident nucleons go, rapidity \sim max.

Relativistic Kinematics @ Collider



$<=$ Fragmentation Central Region Fragmentation $=>$

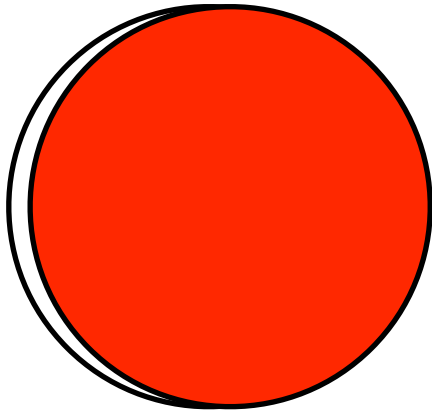
$\# \text{ particles} \uparrow$
(int'd over p_t)



AA collisions: Central vs Peripheral

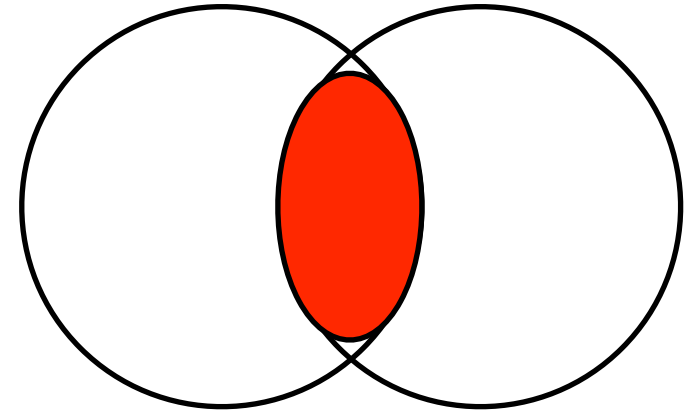
Central:

Maximum
Overlap



Peripheral=>

“Almond” of
overlap region



Theoretically: would like to compare central AA from small to large A.
Takes a lot of beam time. But running with given A, *automatically* measure peripheral collisions.

Exp. variable: # participants.

= 400 in central (= 200 + 200)

= 100 => 400 in peripheral (Glauber & other models; agree to 10%)

Typical Heavy Ion Event @ RHIC

Total # particles = 1000's. Exp.'y: dealing with **high** event rates, data acquisition...
AA @ RHIC similar to pp collisions @ LHC.

Experiments @ RHIC:

STAR: big, 4 π coverage, $y = \pm 2$

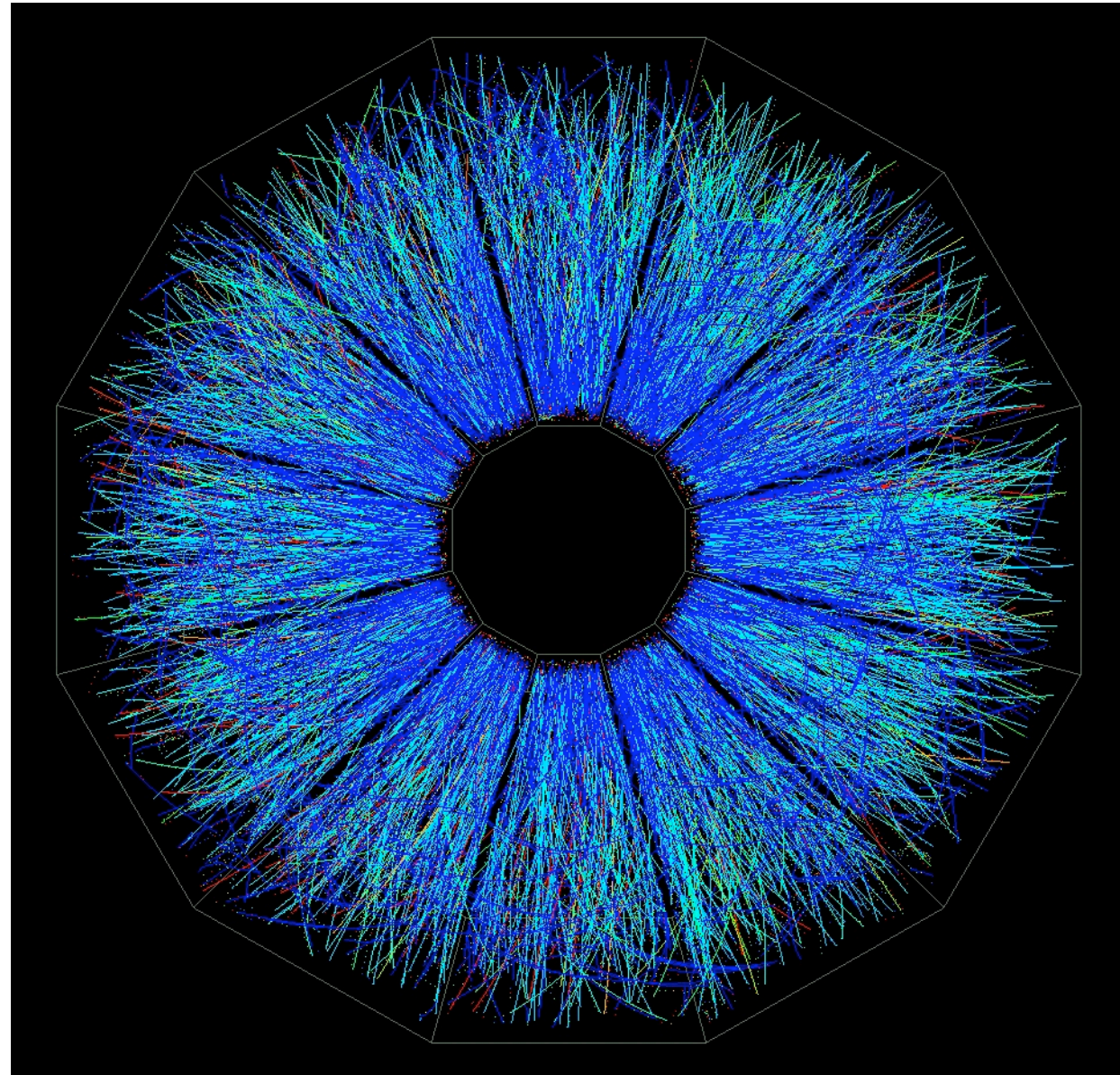
PHENIX: big, elec.-mag., $y = \pm 2$

PHOBOS: small, all rapidity

BRAHMS: small, all rapidity

big = 400 experimentalists
(~ “**participants**”)

small = 50 exp.'s.



The “Body” of the Unicorn: Soft Momenta, $p_t < 2 \text{ GeV}$

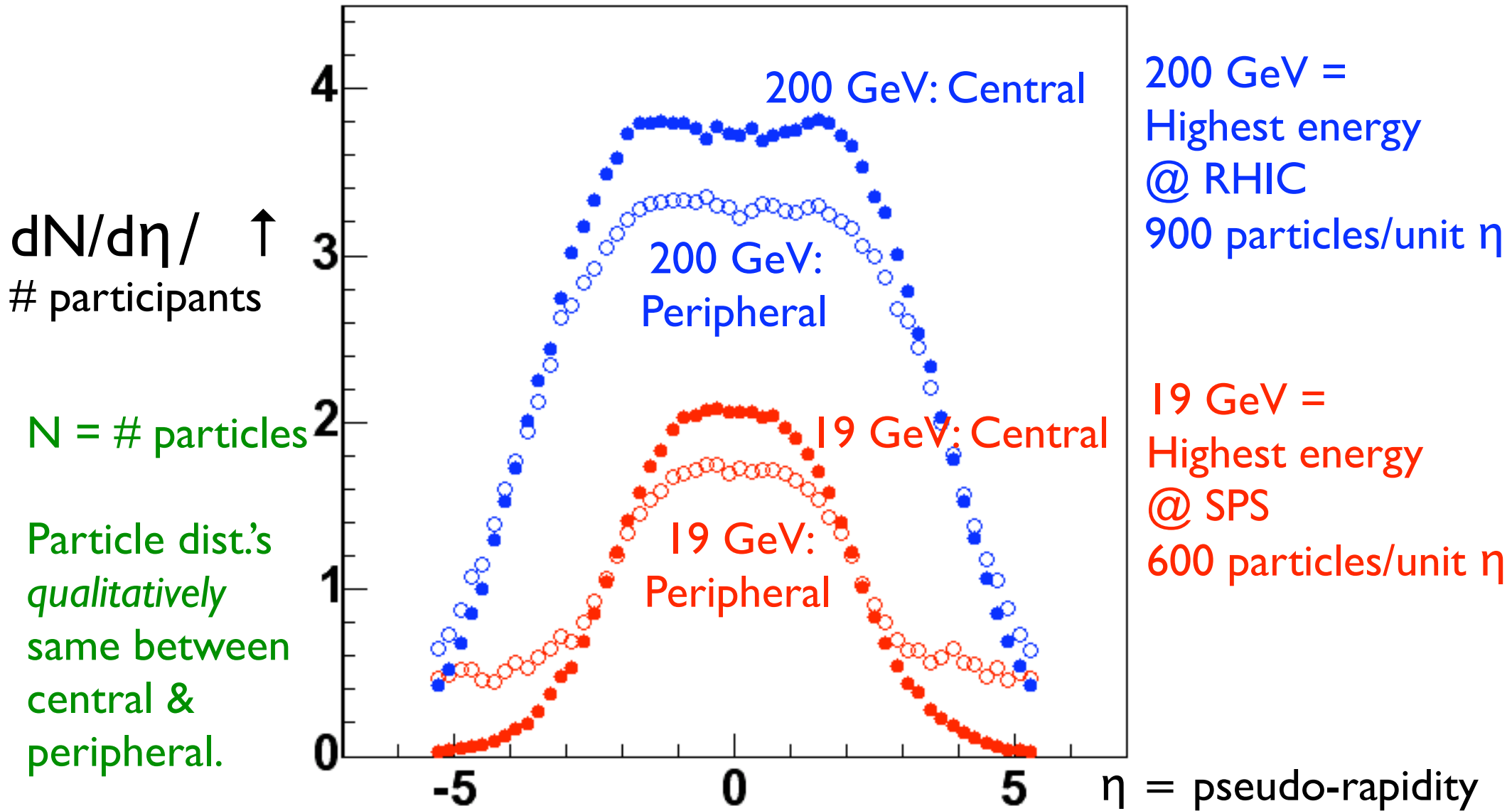
Most particles are at soft momentum.

With $T_c \sim 200 \text{ MeV}$, expect thermal particle distributions to $p_t \sim 2 \text{ GeV}$.

Thousands of particles, should be able to use hydrodynamics...



Particle Distributions vs η , Energy: “Central Plateau” @ RHIC



$\eta \sim 0 = 90^\circ$ for collider. **central region:** $\eta = \pm 2$ @ RHIC

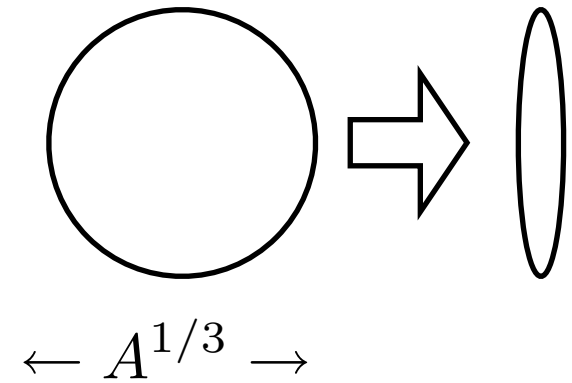
$\eta \sim$ maximum = down the beam pipe. **fragmentation region:** $|\eta| > 2$ @ RHIC

Why do AA? “Saturation” as a Lorentz Boost

At high energies, incident nucleus is *Lorentz contracted*.
=> color charge of incident nucleus gets “squashed”.

McLerran & Venugopalan: color charge bigger by $A^{1/3}$

$A \rightarrow \infty$: can use *semi-classical* methods.



@ central rapidity, *gluon saturation* = **Color Glass**.

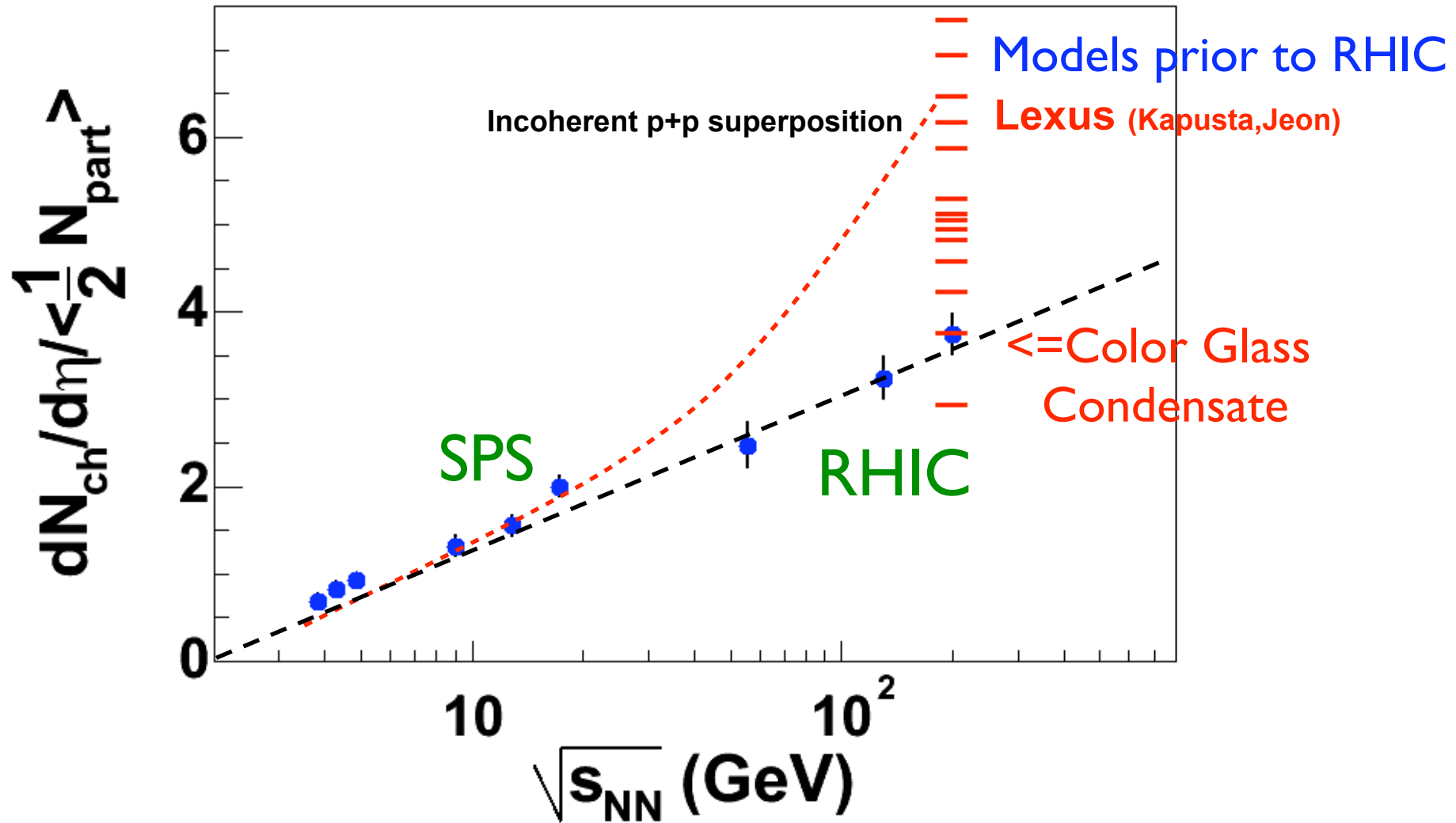
As semi-classical, predicts *logarithmic* growth in multiplicity:

$$\frac{dN}{dy} \sim \frac{1}{g^2(\sqrt{s}/A)} \sim \log(\sqrt{s}/A)$$

First surprise from Day 1: NO big increase in multiplicity. Approx. log growth.

Also: expect avg. momentum to grow similarly $\langle p_t \rangle \sim \log(\sqrt{s}/A)$
(Krasnitz & Venugopalan)

Slow Growth in Multiplicity with Energy

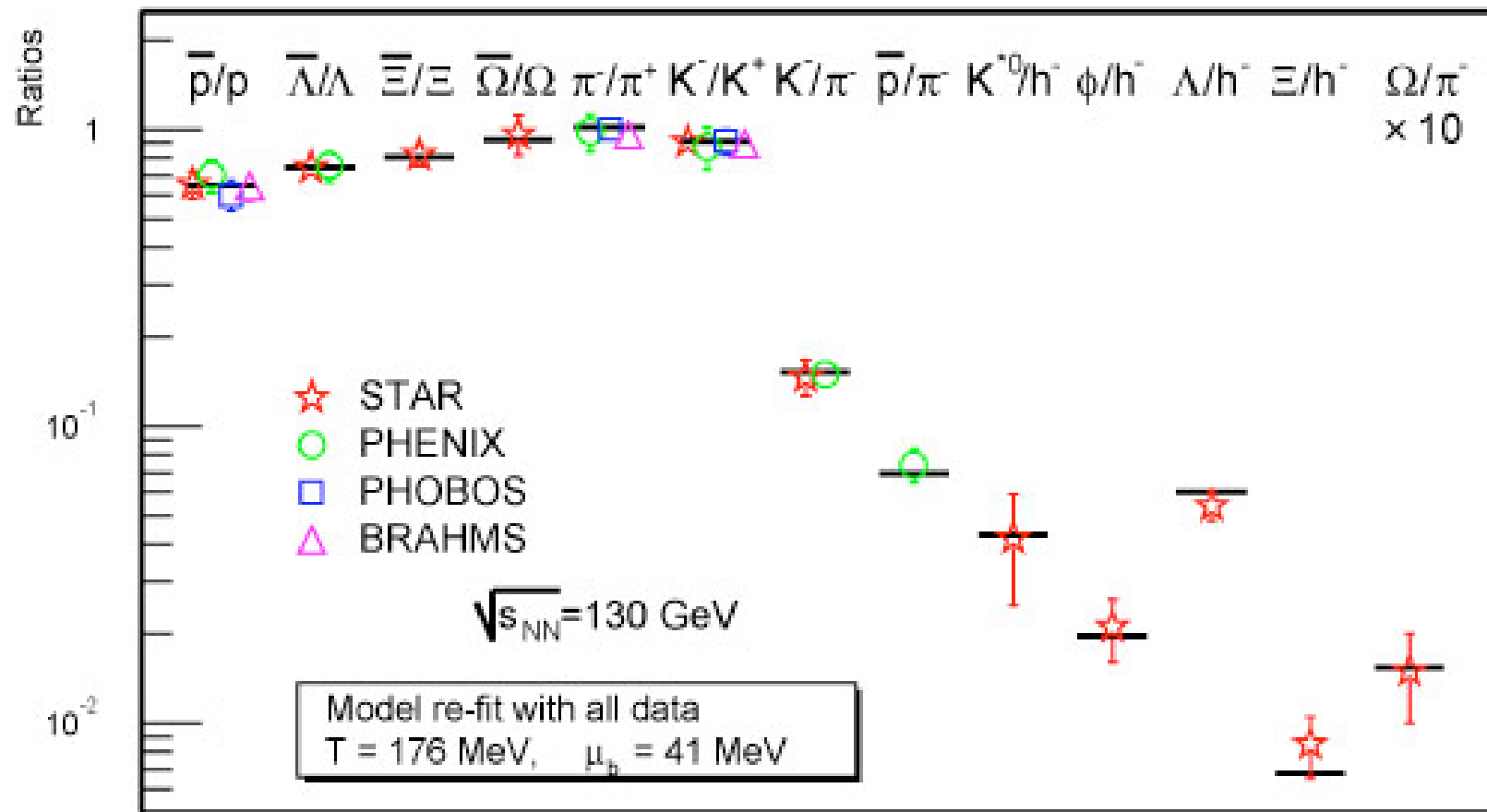


Good fits to overall multiplicity, centrality dependence (Kharzeev, Levin, Nardi)

But: STAR: from 130 \Rightarrow 200 GeV, multiplicity increases by 14%,
but NO change in $\langle p_t \rangle \pm 2\%$. Vs. $> 7\%$ increase from Color Glass!

Total Chemical Ratios *Appear* in Thermal Equilibrium

$$T_{ch} = 175 \text{ MeV}$$



Braun-Munzinger et al., PLB 518 (2001) 41 D. Magestro (updated July 22, 2002)

OVERALL chemical abundances *well* fit with $T_{ch} = 175 \text{ MeV}$, $\mu_{\text{baryon}} \sim 0$
 (Becattini, Braun-Munzinger, Letessier, Rafelski, Redlich, Stachel, Tounsi...)

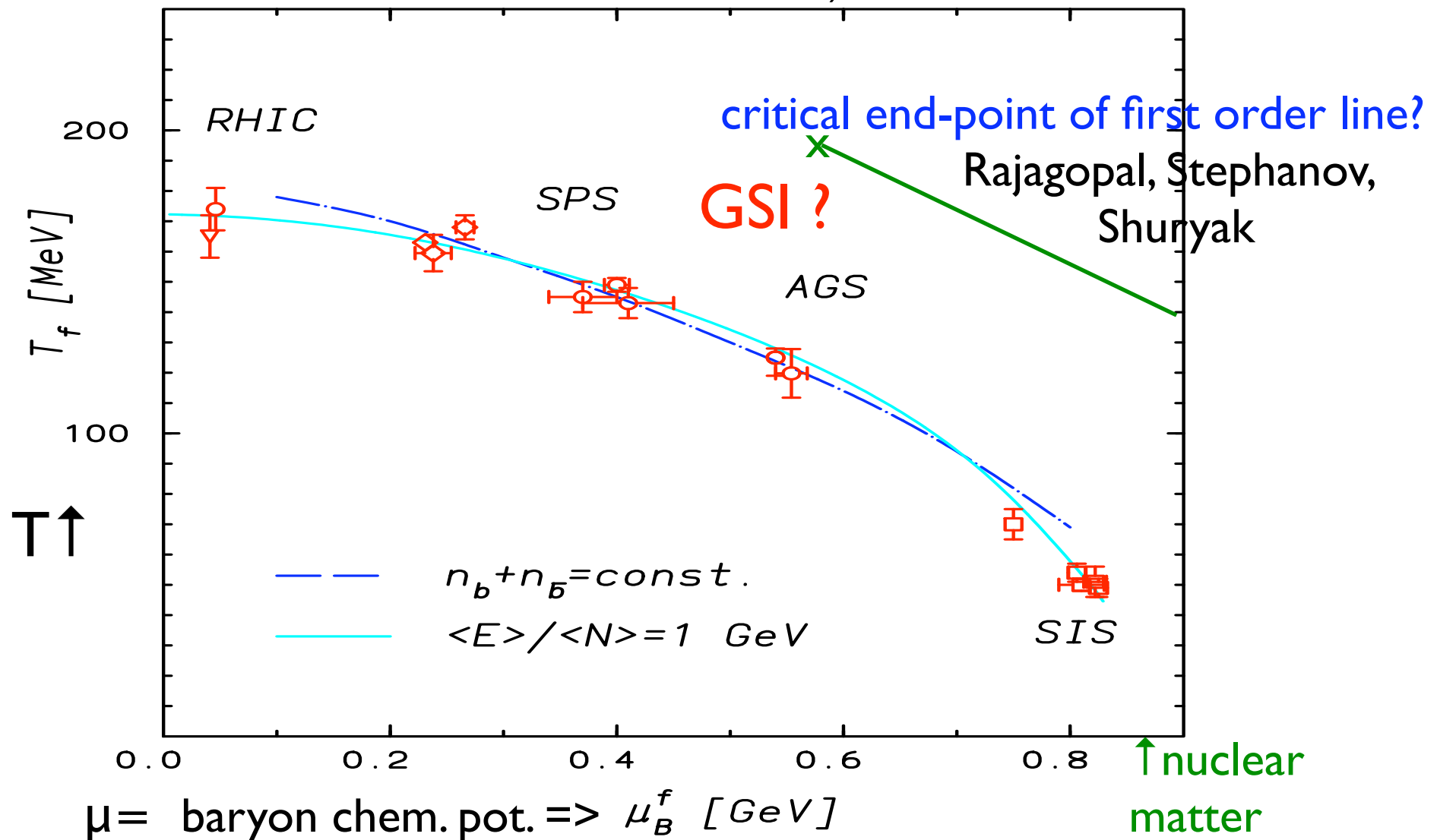
N.B.: even for multi-strange baryons, with relative abundances $\sim 1\%$ of pions.

Chemical Ratios vs Energy in AA: T- μ plane

Similar fits for chemical abundances also work at lower energies. Baryons still present at $y=0$, so need to add **baryon chemical potential, μ** .

Find *line* in T- μ plane. **Similar fits work for pA, pp - everywhere!**

(With corr.'s for finite vol., canonical ensemble...) **=> NOT conclusive.**



p_t Spectra *Appear* In Thermal Equi. \sim *Hydrodynamics*

$T_{kin} \approx 100 \text{ MeV} (\ll T_{ch}!)$ Local Boost Velocity $\beta \sim .7c$

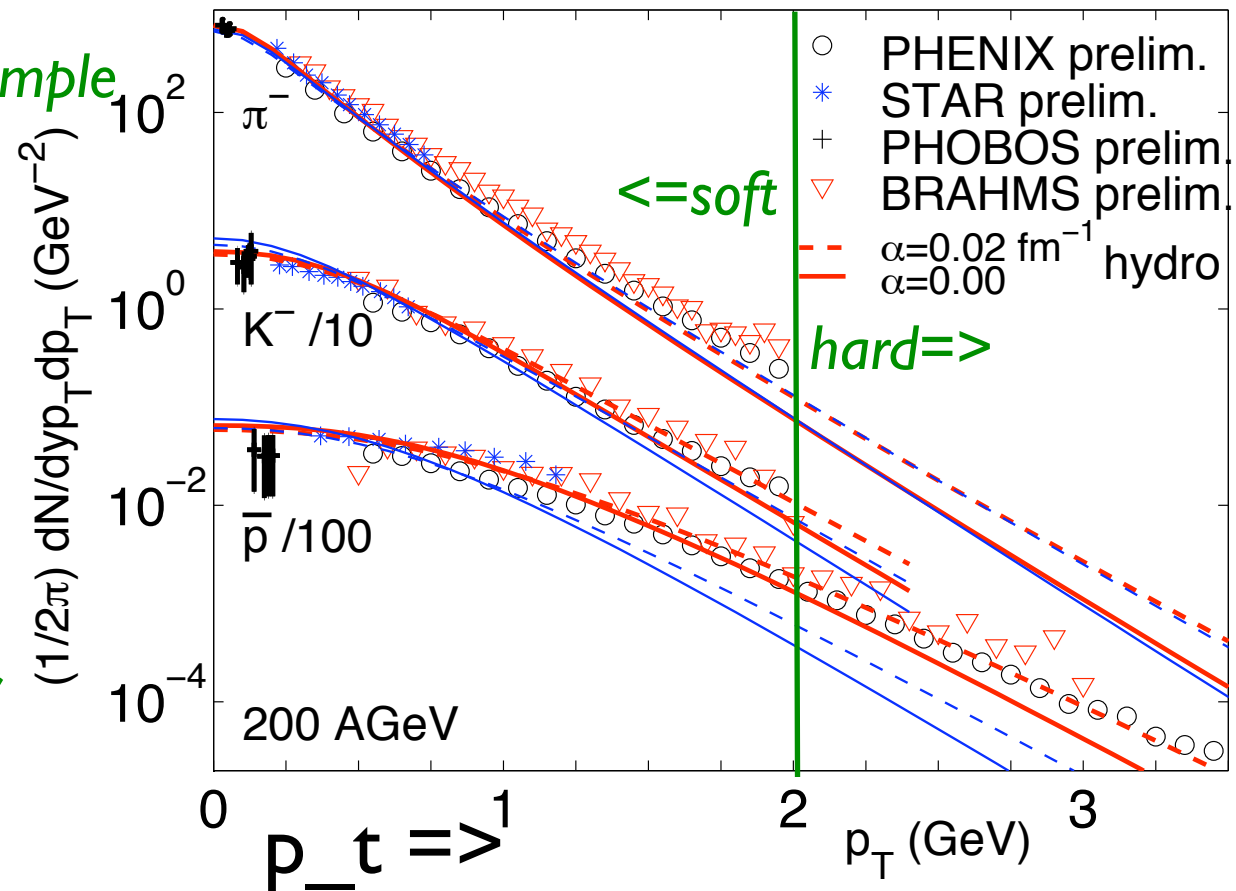
Hydro. gives good description for most particles, at low $p_t < 1 \text{ GeV}$.

Assumes initial conditions: starts
above T_c in thermal equilibrium, *simple*
Equation of State (1st order!)
Ideal hydro.: NO viscosity...

Large local boost velocity $\beta \sim .7 c$.
Spectra of heavy particles “turn
over” at low p_t . $\beta = \beta(\text{radius})$.

RHIC: first clear evidence for
boost velocity: big!

Direct fits similar: “Blast-wave”

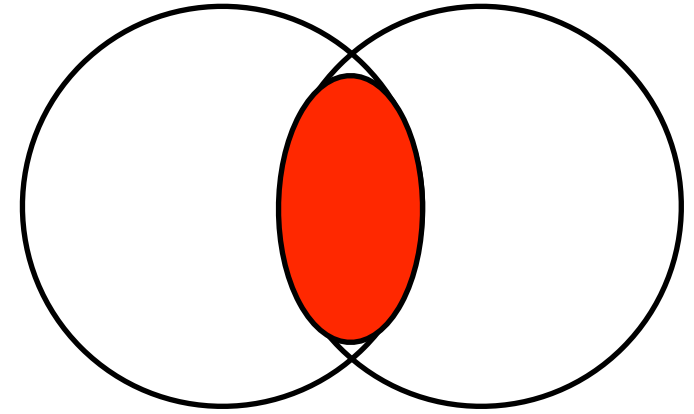


Hydro *needs* to assume applicable from very early times, $.6 \text{ fm}/c$!

Heinz, Hirano, Kolb, Rapp, Shuryak, Teaney... (above Heinz & Kolb)

Success of Hydro.: $v_2 = \text{Elliptical Flow}$

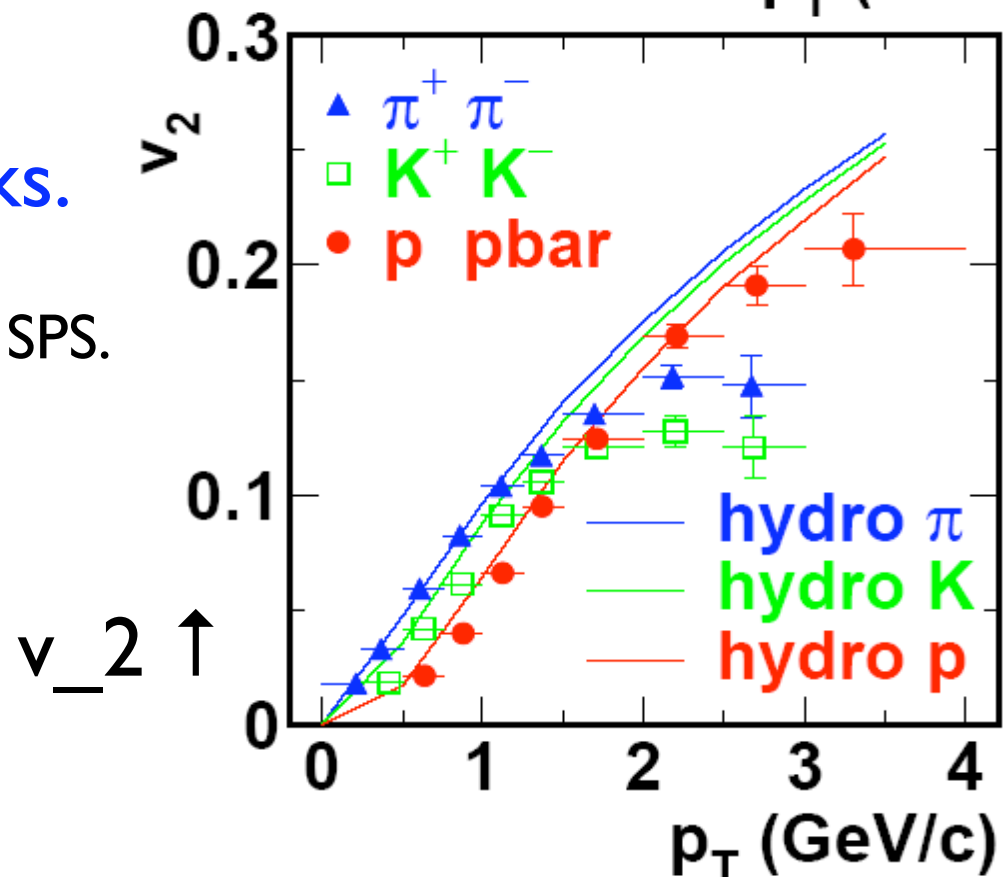
Peripheral Coll.'s: Start with system which is anisotropic in momentum space. Exp.'y, compute how *spatial* anisotropy \Rightarrow *momentum* anisotropy. (Ollitrault, Borghini)



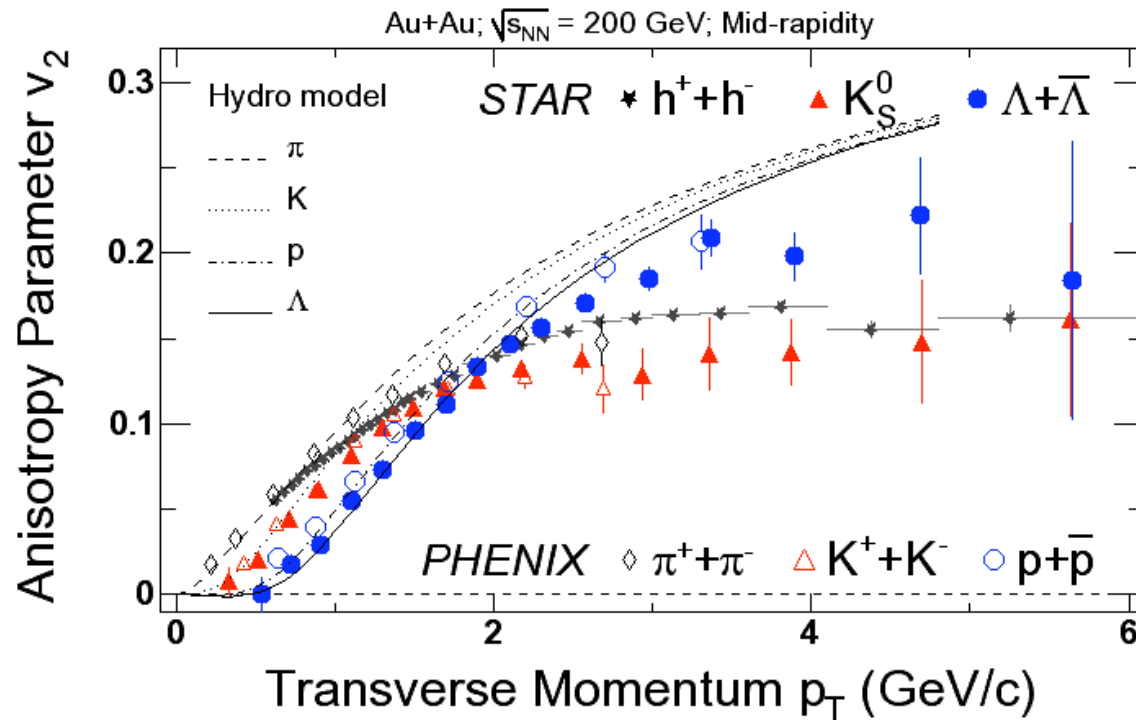
$$v_2 = \langle \cos(2\phi) \rangle, \quad \tan \phi = p_y / p_x$$

$v_2 \Rightarrow$ collective behavior:
there is “stuff”, and it sticks.

Hydro works for v_2 @ RHIC, not SPS.



At Low $p_t < 1$ GeV, Hydro. works for All Particles

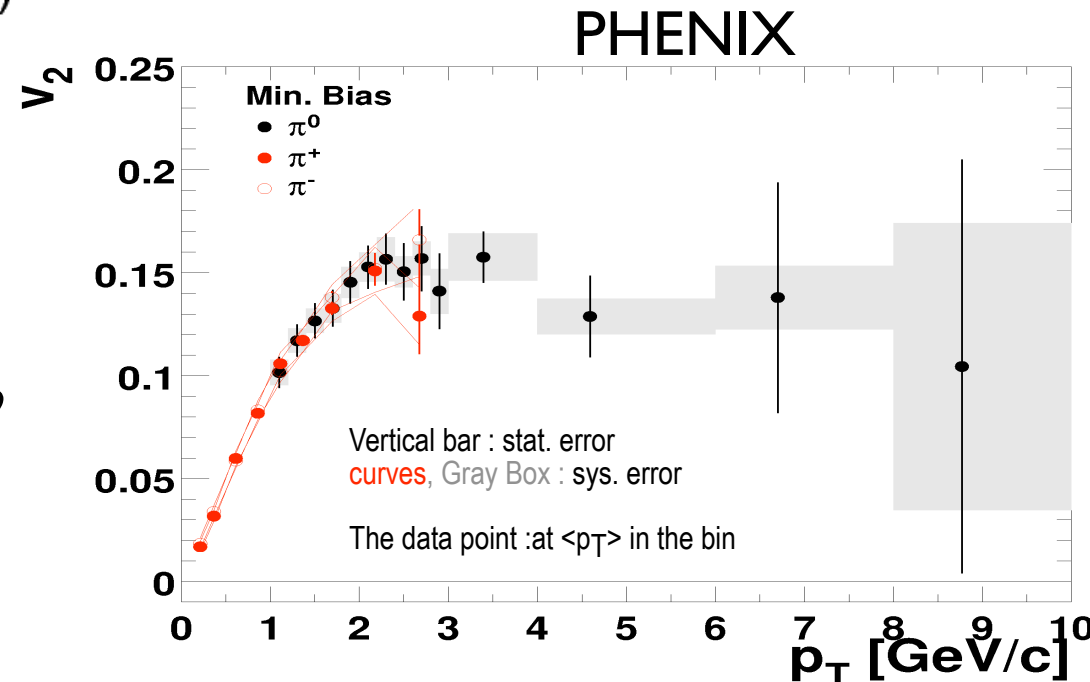


\leq Hydro works for v_2
to $p_t \sim 1$ GeV for
 π 's, K 's, p 's, Λ 's.... everything.

For all particles, v_2 flat for
 $p_t > 1$ GeV \Rightarrow 10 GeV - !!

Is v_2 at $p_t > 1$ GeV measuring
collective flow, or jet-jet correlations?
Apparently: true collective flow.

So why flat?



HBT Radii: Hydro *Fails*. “Blast Wave” Works

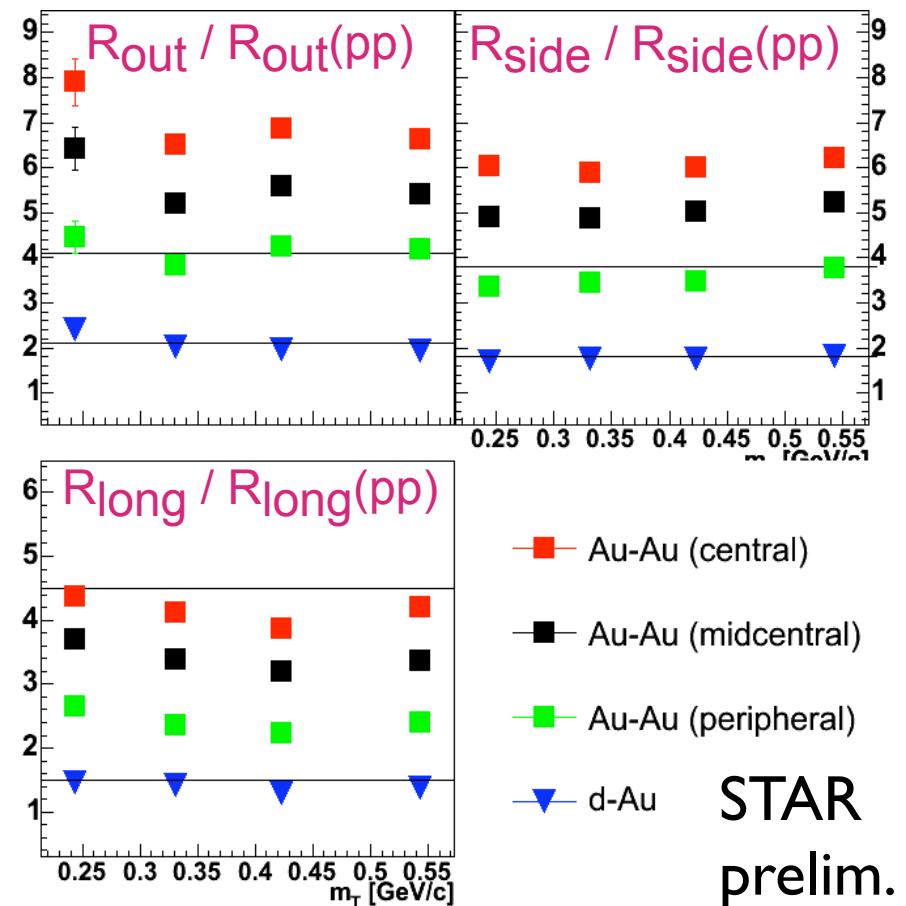
Hanbury-Brown-Twiss (HBT) radii: two-particle correlations for identical particles, used to determine size (as for stars). Typically: fall off like Gaussian.

Here: *three* directions for momentum of pion pair (Bertsch & Pratt).
HBT then gives three sizes: along beam ($R_{\text{longitudinal}}$), along line of sight (R_{out}), & perpendicular to light of sight (R_{side}).

Hydro: $R_{\text{out}}/R_{\text{side}} > 1$, *increases* with p_{t} .
Exp.: $R_{\text{out}}/R_{\text{side}} \sim 1$, *decreases* with p_{t} !

“Blast Wave” works: expanding shell.
Is a *fit*, not underlying space-time picture.

HBT radii \sim *same* in pp, dA, and AA!
Even p_{t} dependence same!



Body of the “Unicorn”:

Majority of particles, at small momenta < 2 GeV, look superficially like thermal bath. But in detail, surprises:

HBT radii =>
space-time picture *not* yet understood.

Tail of the “Unicorn”:

Look at particles at *HIGH* momentum, $p_t > 2$ GeV, to probe the body.

The Tail wags the (Dog) Unicorn

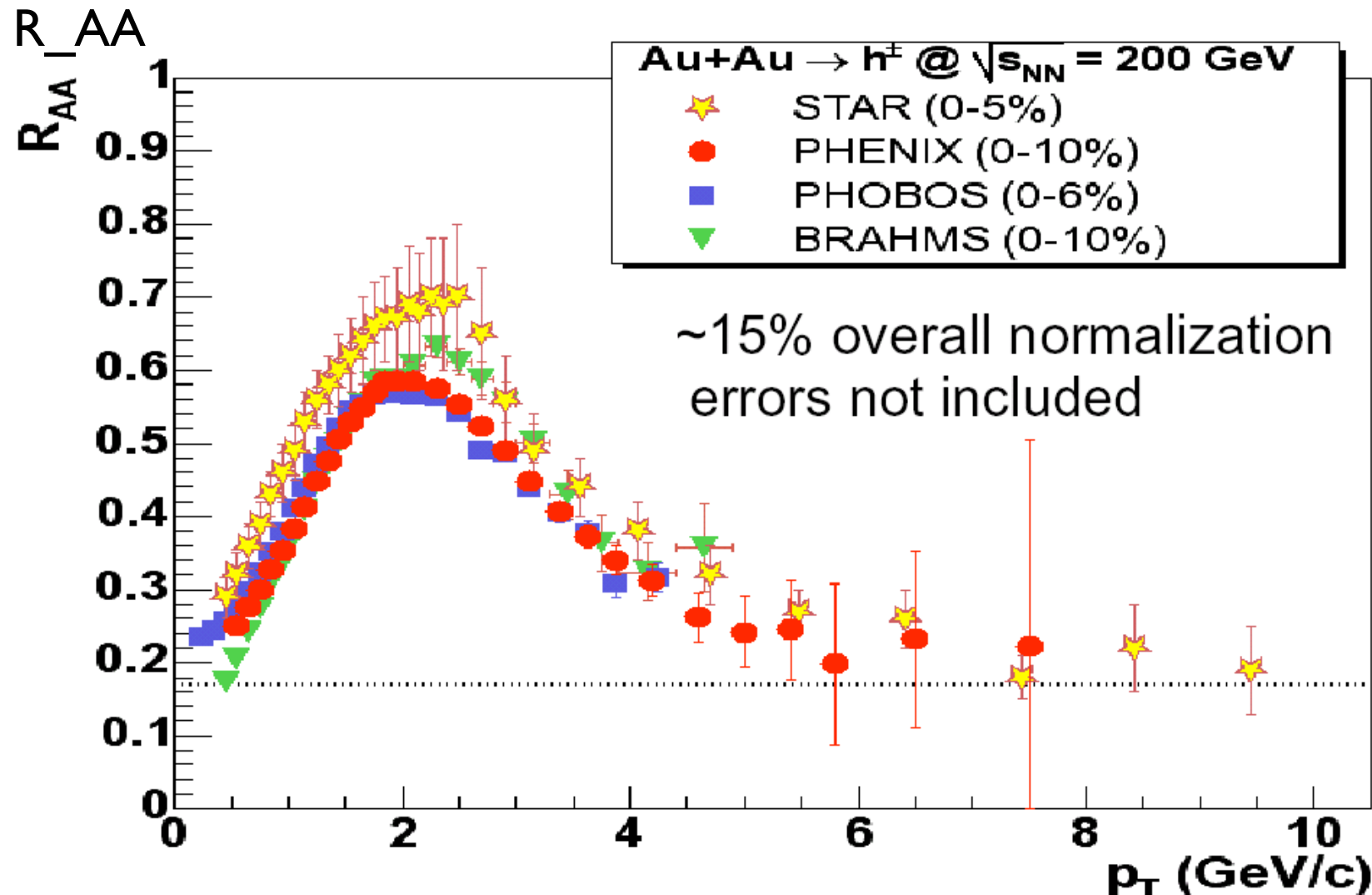


Clear Experimental Signal of “Stuff”: R_{AA}

Compare spectra in AA to that in pp, especially for “hard” $p_t > 2$ GeV:

From Day 1, “hard” spectra appear *steeper* in AA than pp \Rightarrow fewer particles.

$R_{AA} = \frac{\# \text{ particles at a given } p_t \text{ in central AA collision}}{\# \text{ particles at the same } p_t \text{ in pp, central rapidity.}}$



$R_{AA} \Rightarrow$
*suppression of
hard particles
in AA, vs pp.*

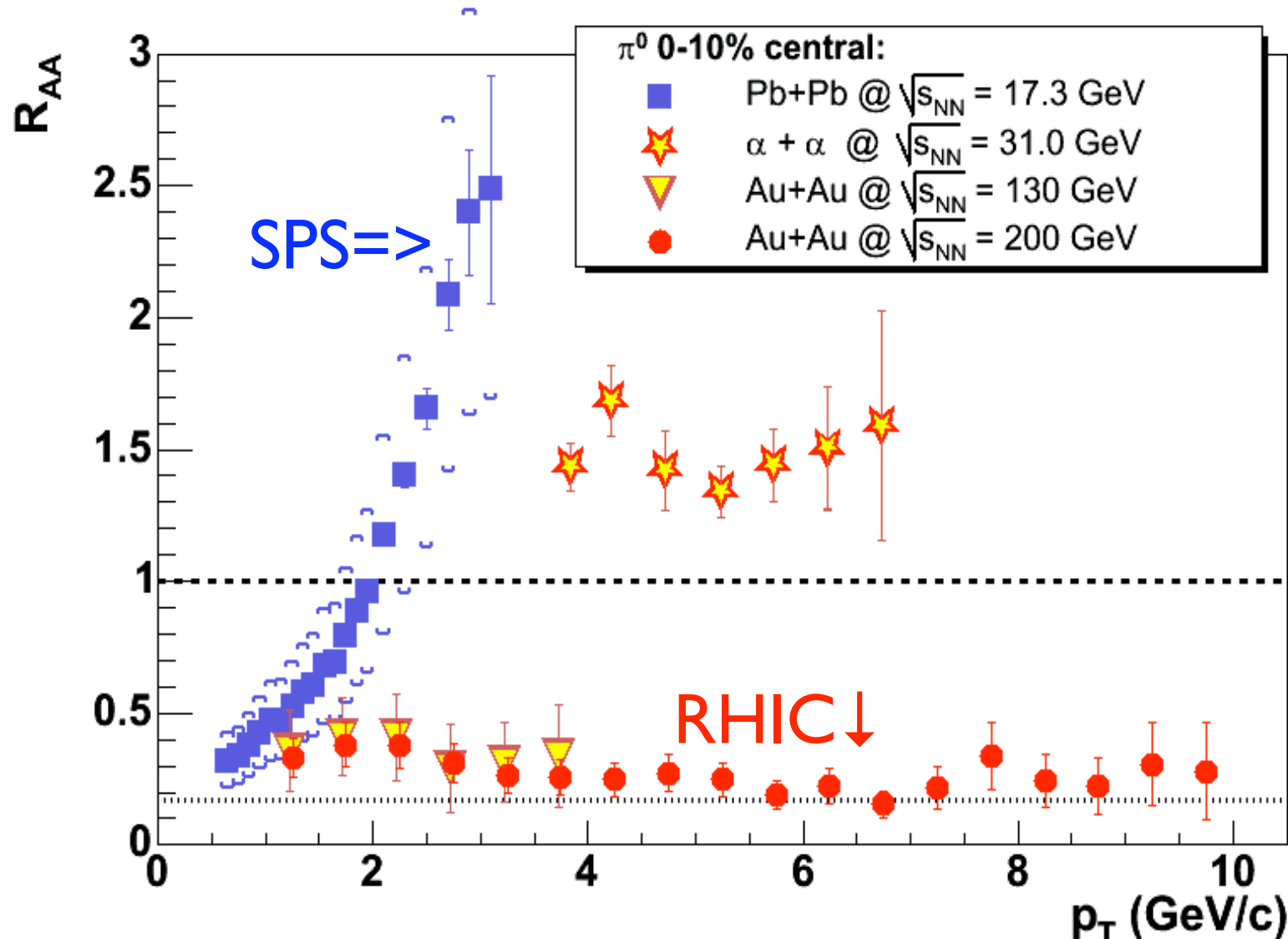
*For $p_t > 6$
GeV
all particles
suppressed.*

R_AA: Enhancement @ SPS, Suppression @ RHIC

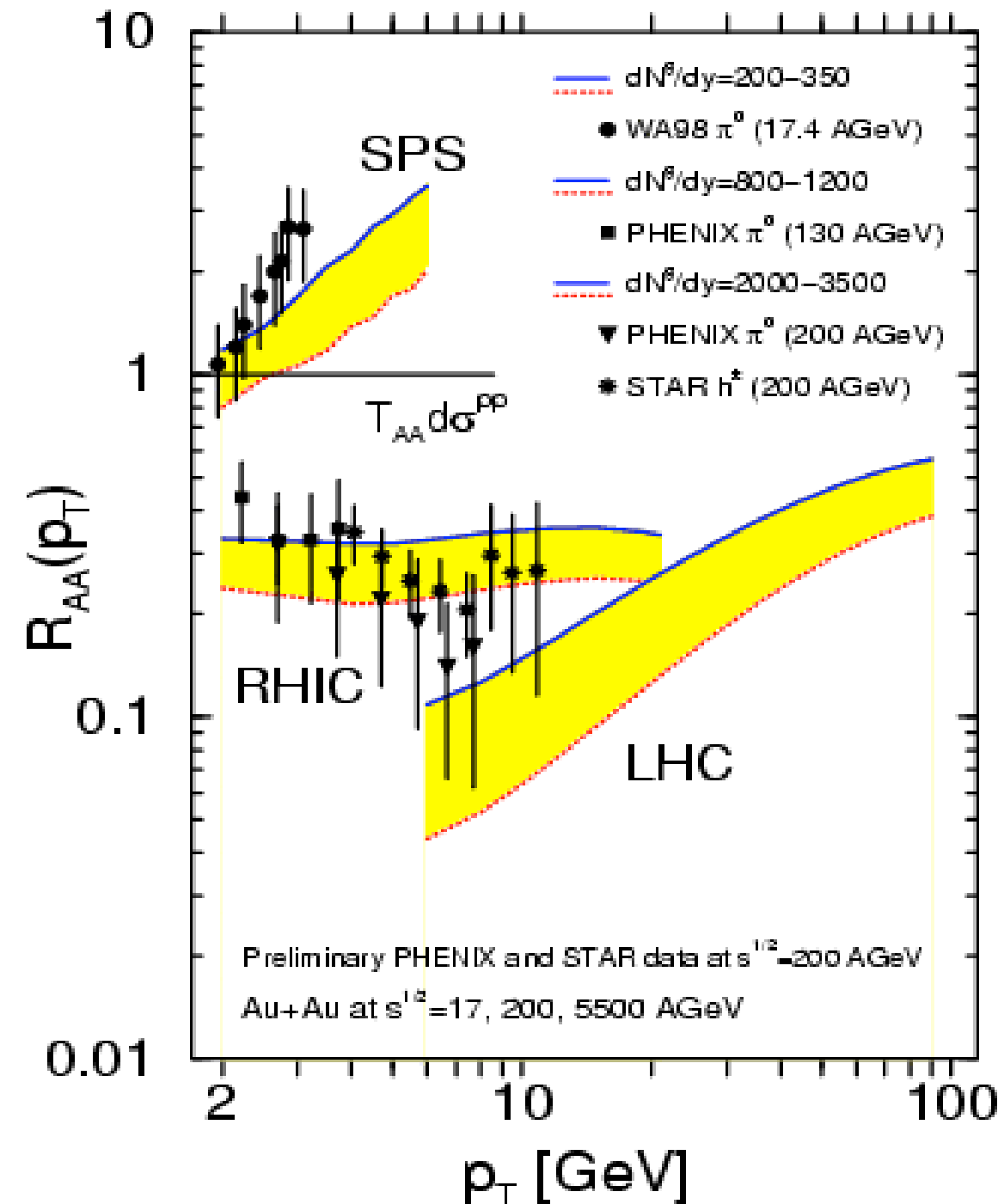
Effect most dramatic for π^0 's. SPS: $R_{AA} \sim 2.5$ @ 3 GeV. “Cronin”

RHIC: $R_{AA} \sim 0.2$ @ 3 GeV.

RHIC: Supp. from energy loss - “stuff” slows fast particles down.



R_AA: Qualitative Agreement with “Energy Loss”



Energy Loss: A fast particle going through a thermal bath loses energy:

Gyulassy, X.N. Wang, Vitev...Baier, Dokshitzer, Mueller, Schiff, Zakharov

\leq Gyulassy & Vitev: *conspiracy* to give *flat* R_{AA} @ RHIC.

Need to add several effects, “Cronin”, energy loss, shadowing...

Is “flat” R_{AA} for π^0 ’s special to RHIC? Will be interesting @ LHC!

Central AA: at inter. p_t , *only* mesons suppressed

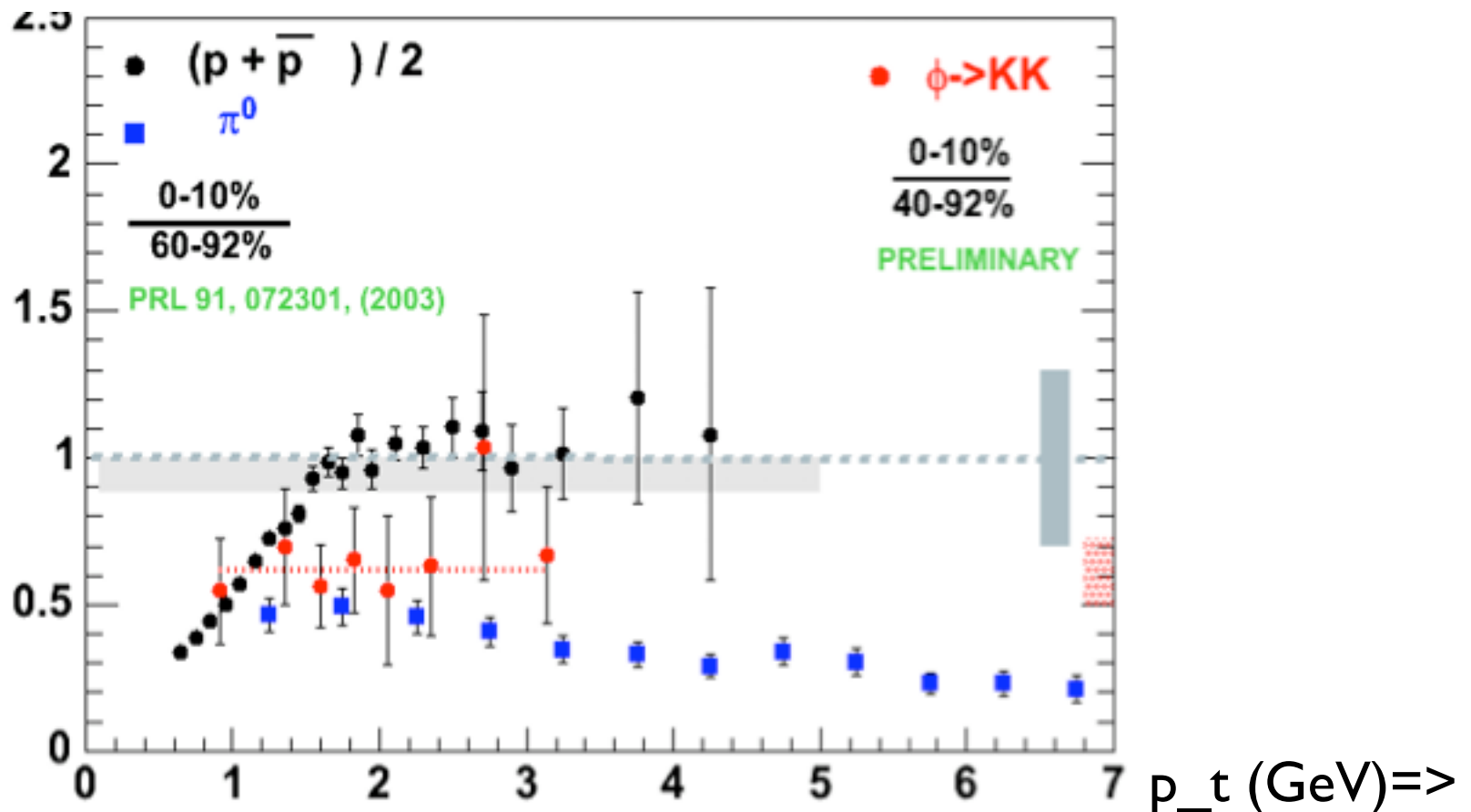
R_{CP} : ratio for # particles at given p_t , for **central to peripheral collisions**

Behaves like R_{AA} , easier to get data.

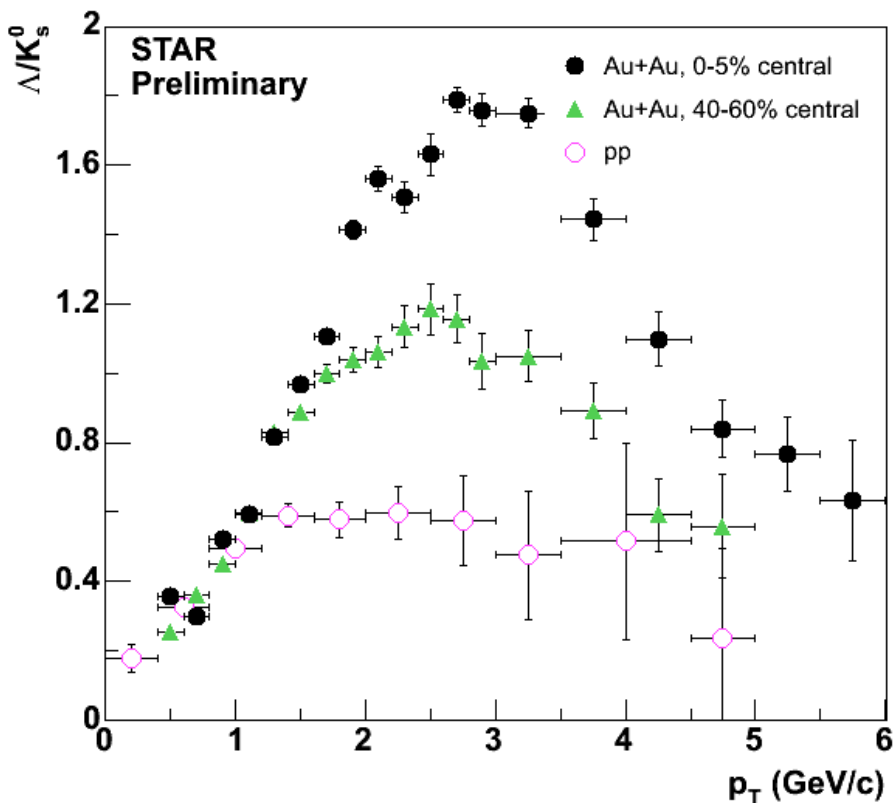
Find: *baryons* not suppressed for $p_t: 2 \Rightarrow 6$ GeV, *mesons* are.

Mesons suppressed \Rightarrow “stuff” is gluonic.

$R_{CP} \uparrow$



Baryon “Bump” at $p_T: 2 \Rightarrow 6$ GeV



Central AA: *baryon “bump” at $p_T: 2 \Rightarrow 6$ GeV*

Baryon/meson ratio enhanced by ~ 3 in central AA vs pp. First seen in p/π .

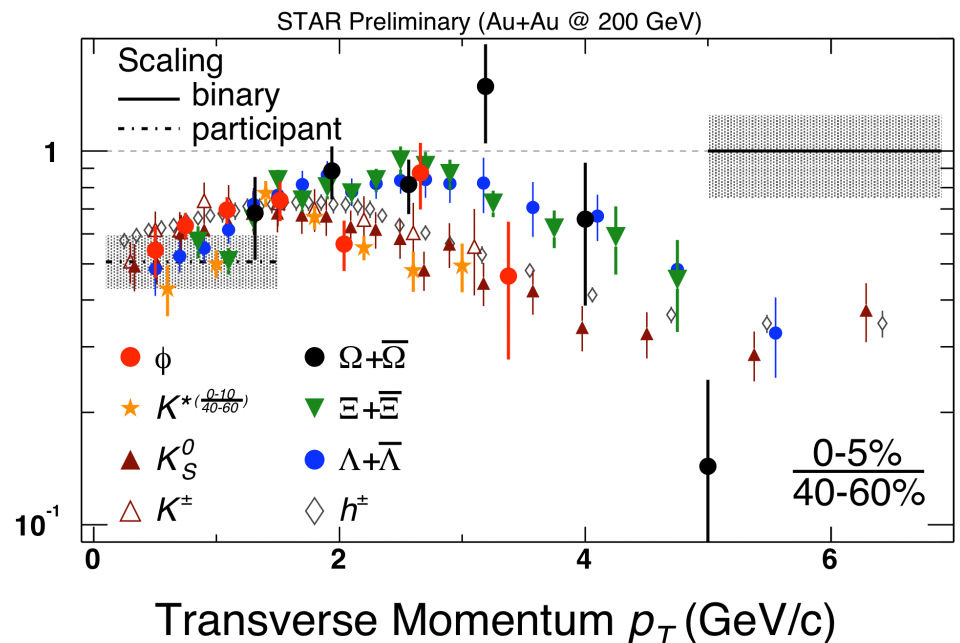
$\leq \Lambda/K$ ratio: bump peaks at ~ 3 GeV.

Above $p_T = 6$ GeV, ratios like pp.

R_{CP} vs particle species \Rightarrow

All particles suppressed > 6 GeV, $R_{CP} \sim 0.2$.

\Rightarrow Gluon “stuff” supp.’s mesons, generates baryon “bump”

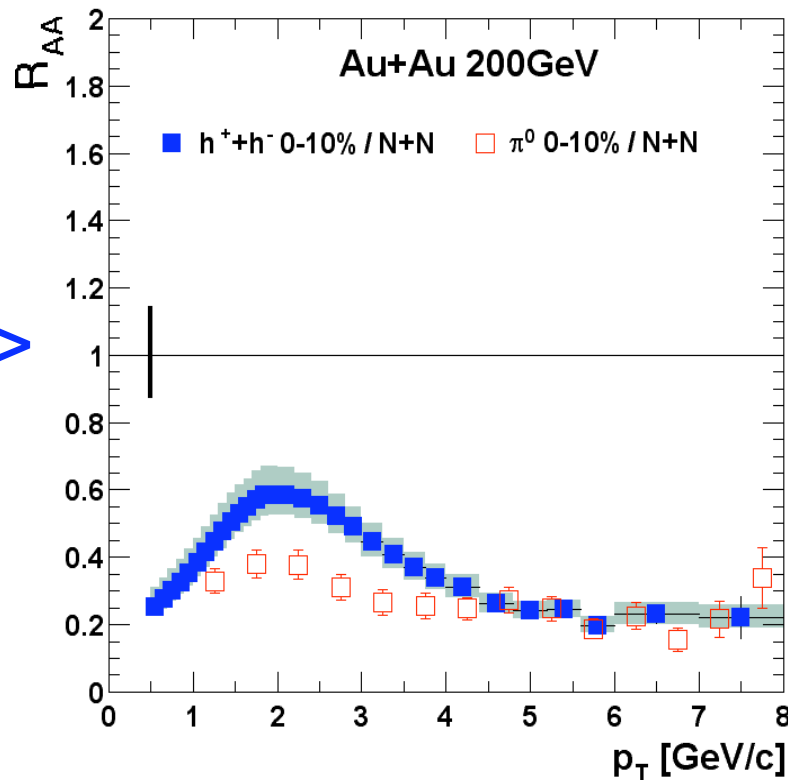


R_{AA} Final State Effect: *NOT* seen in R_{dA}

Look at R_{dA} , analogous ratio in dA collisions @ *central rapidity* ($y=0$):
find “Cronin” enhancement in dA, vs suppression in AA.

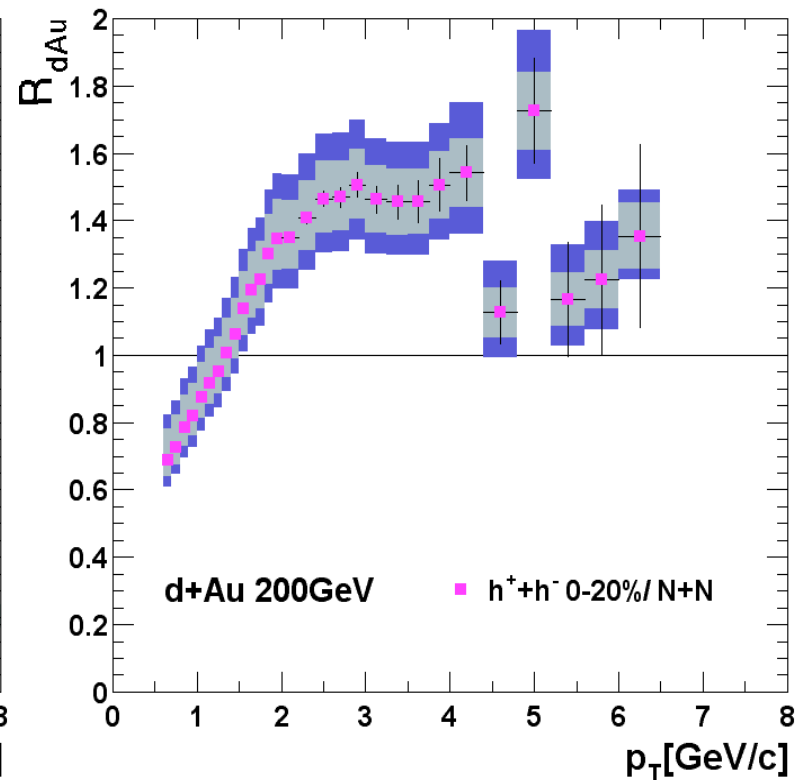
Color Glass (initial state effect) predicted suppression in dA, *not* seen.

AA=>



Suppression in AA \uparrow
 $R_{AA} \sim 0.4$ @ 3 GeV

<=dA



Enhancement in dA \uparrow
 $R_{dA} \sim 1.4$ @ 3 GeV

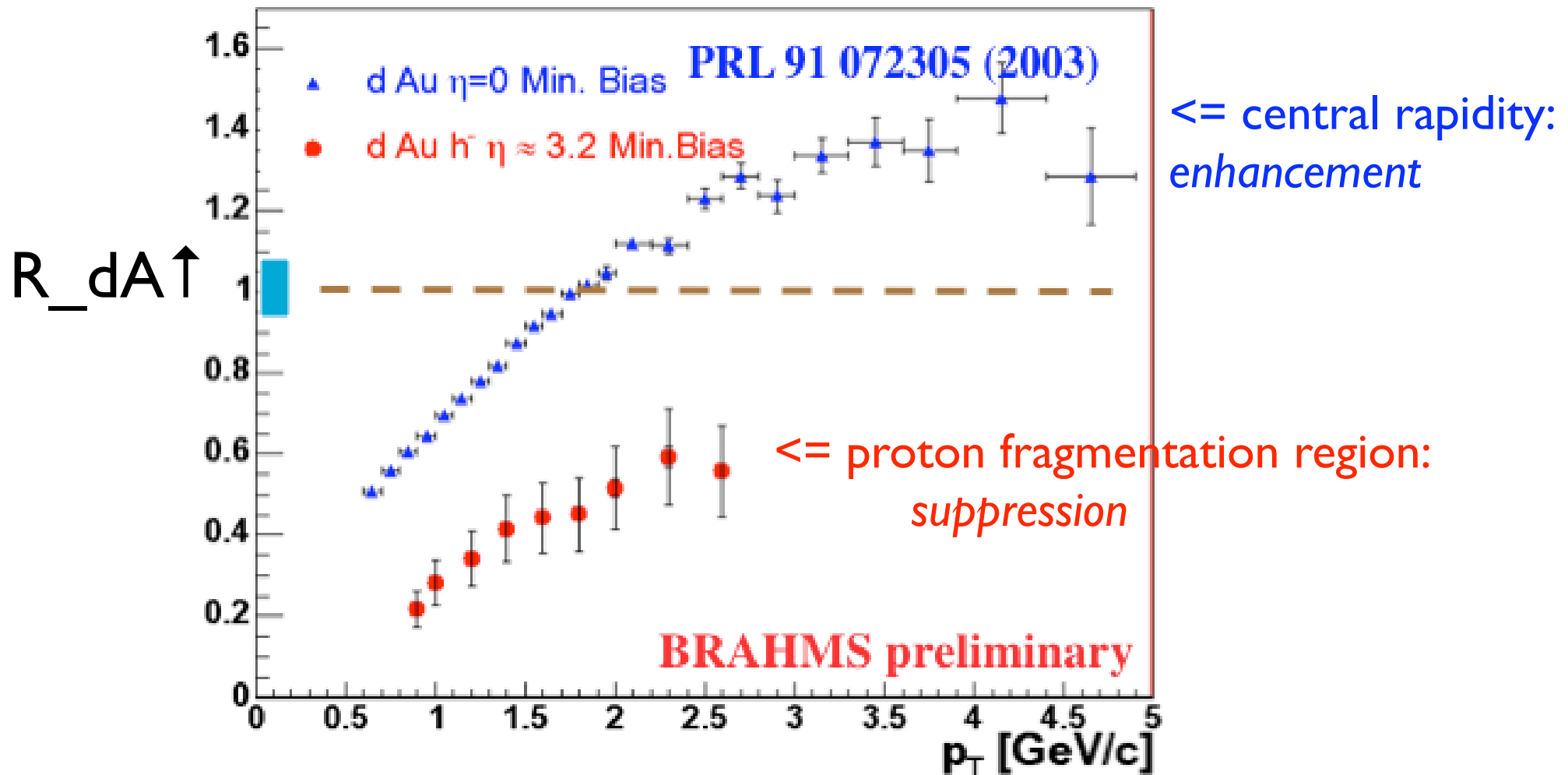
Where to find the Color Glass: dA, by the *proton*

dA: fragmentation region of nucleus tells one about *final* state effects.

frag. region of proton: in the proton rest frame, feels the large color charge of the incident nucleus => *sensitive to initial state effects*:

= *the place to find the Color Glass* (Dumitru, Gelis, Jalilian-Marian)

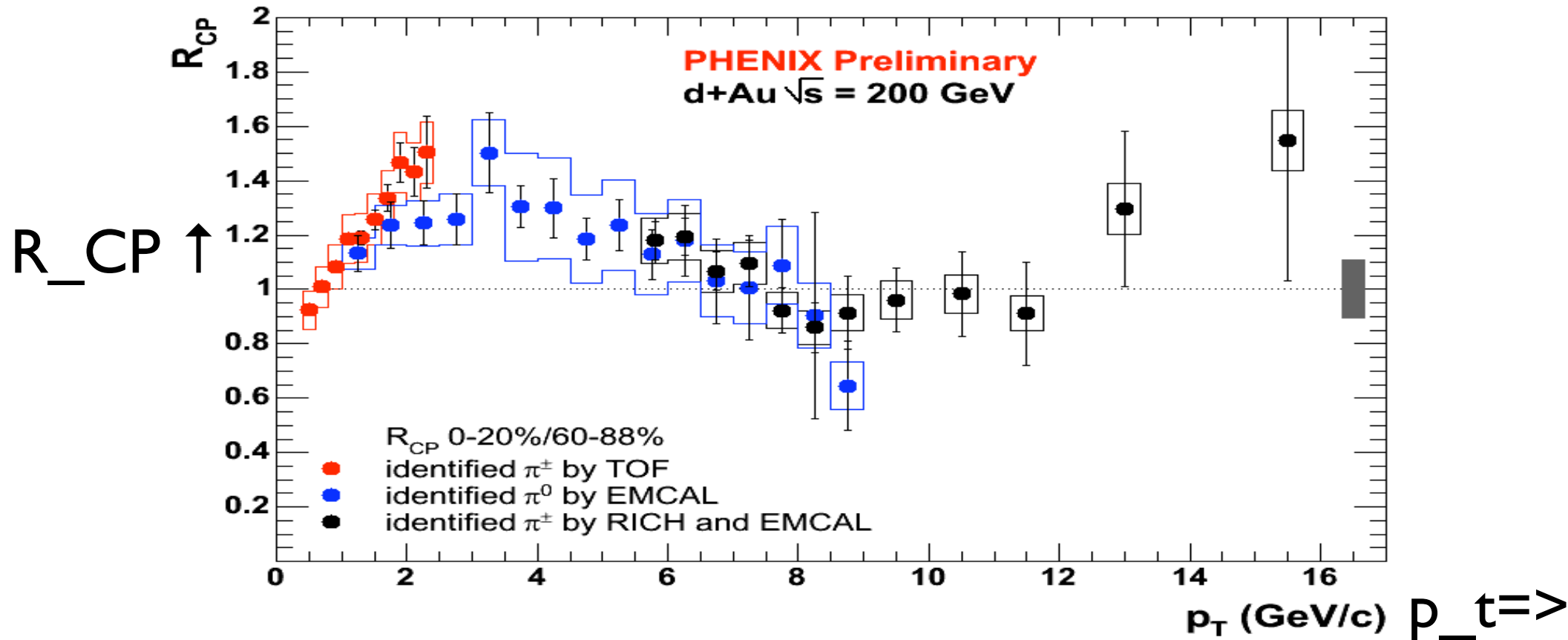
BRAHMS: in dA, *enhancement @ $y=0$, suppression @ proton frag. region.*



dA: No “Cronin” Enhancement at High p_t

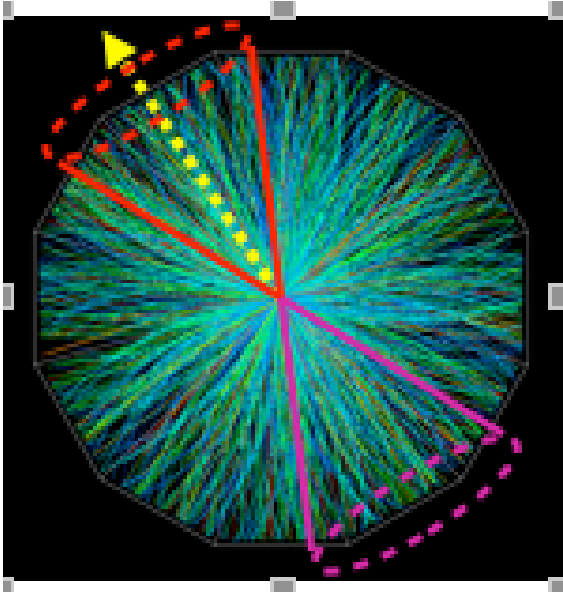
At high p_t , all R's (R_{AA} & R_{CP}) should go to one.

In dA, seen in R_{CP} for $p_t \sim 8$ GeV.



At what p_t does $R_{AA} \Rightarrow 1$? > 10 GeV!

The “Tail” of the Unicorn: Central AA “Eats” Jets



In pp collisions at $\sqrt{s} = 130, 200$ GeV, clearly see “jets”: high energy quarks (& gluons) in each event.

\leq “jet” in AA: cannot see on an event by event basis.

In AA, construct statistical measure: trigger on *hard* particle in one direction, look for *associated* particle in the backward direction

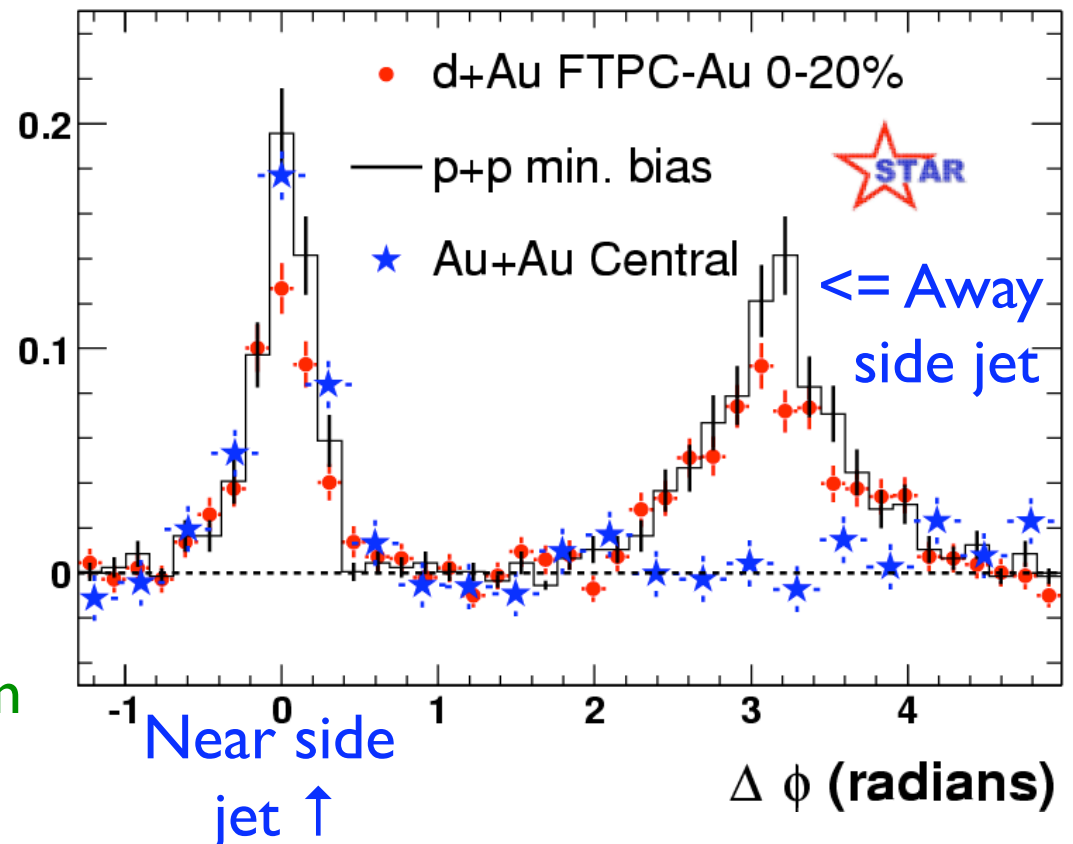
forward: $6 > p_t > 4$ GeV Adams *et al.*, Phys. Rev. Let. 91 (2003)

back: $p_t > 2$

In pp & dA, *clearly* see “backward” peak in angular correlation \Rightarrow associated jet.

In central AA, backward peak is *gone*: “stuff” in AA “eats” jets.

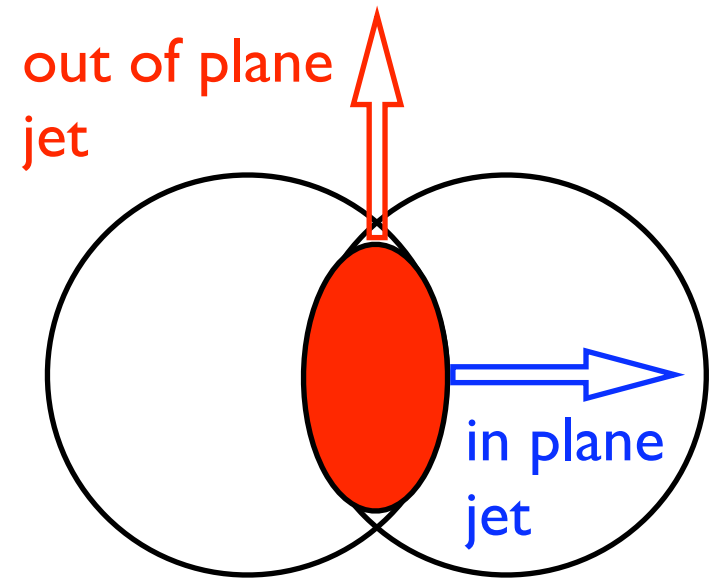
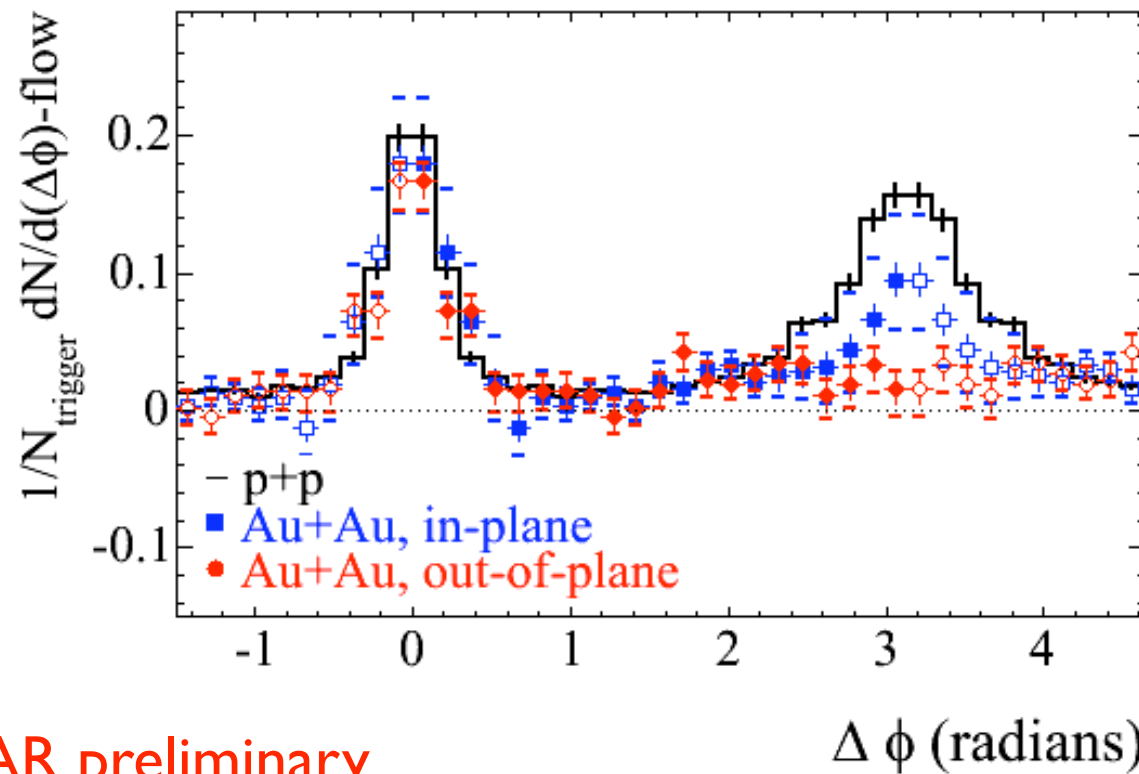
Central AA really “eats” the jet: essentially *nothing* at *hard* momentum in the backward direction.



Peripheral Coll's: Geometrical Test that AA Eats Jets

In peripheral collisions, “stuff” forms an “almond”; a jet has to travel farther through the almond, **out** of the reaction plane, than **in** the reaction plane.

=> Geometrical test that AA “eats” jets: backward jet more strongly suppressed out of plane than in plane!



peripheral collision ↑
almond = “stuff”

STAR preliminary

Suppression larger out-of-plane

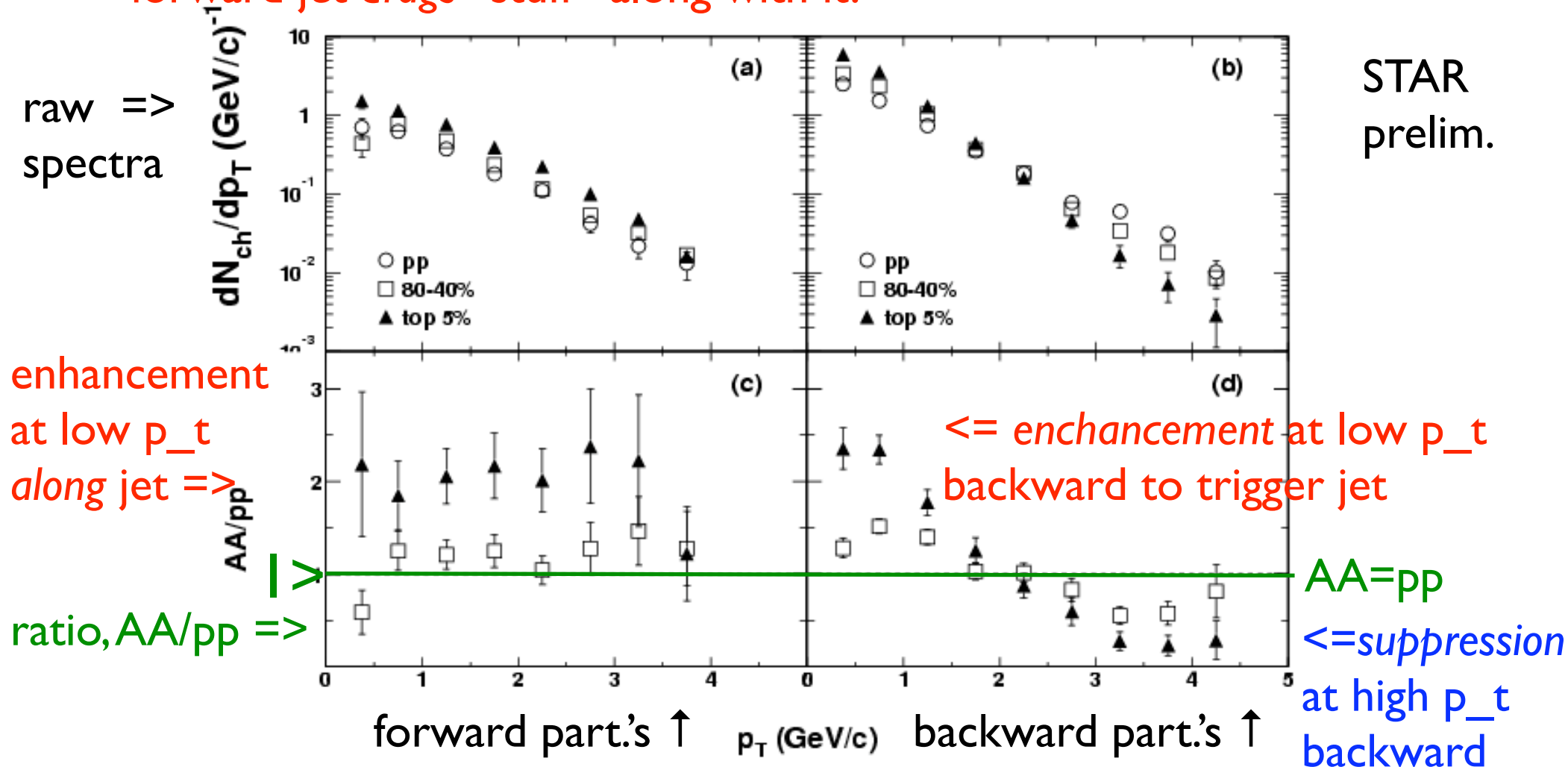
Where does the Backward Jet go in AA?

As before, trigger on forward jet, $6 > p_{t,jet} > 4$ GeV. But look at *all* particles, $p_{t,jet} > .15$ GeV, in both forward and backward directions.

In direction opposite to jet, *suppressed* at high $p_{t,jet}$ (yes), & *enhanced* at low $p_{t,jet}$.

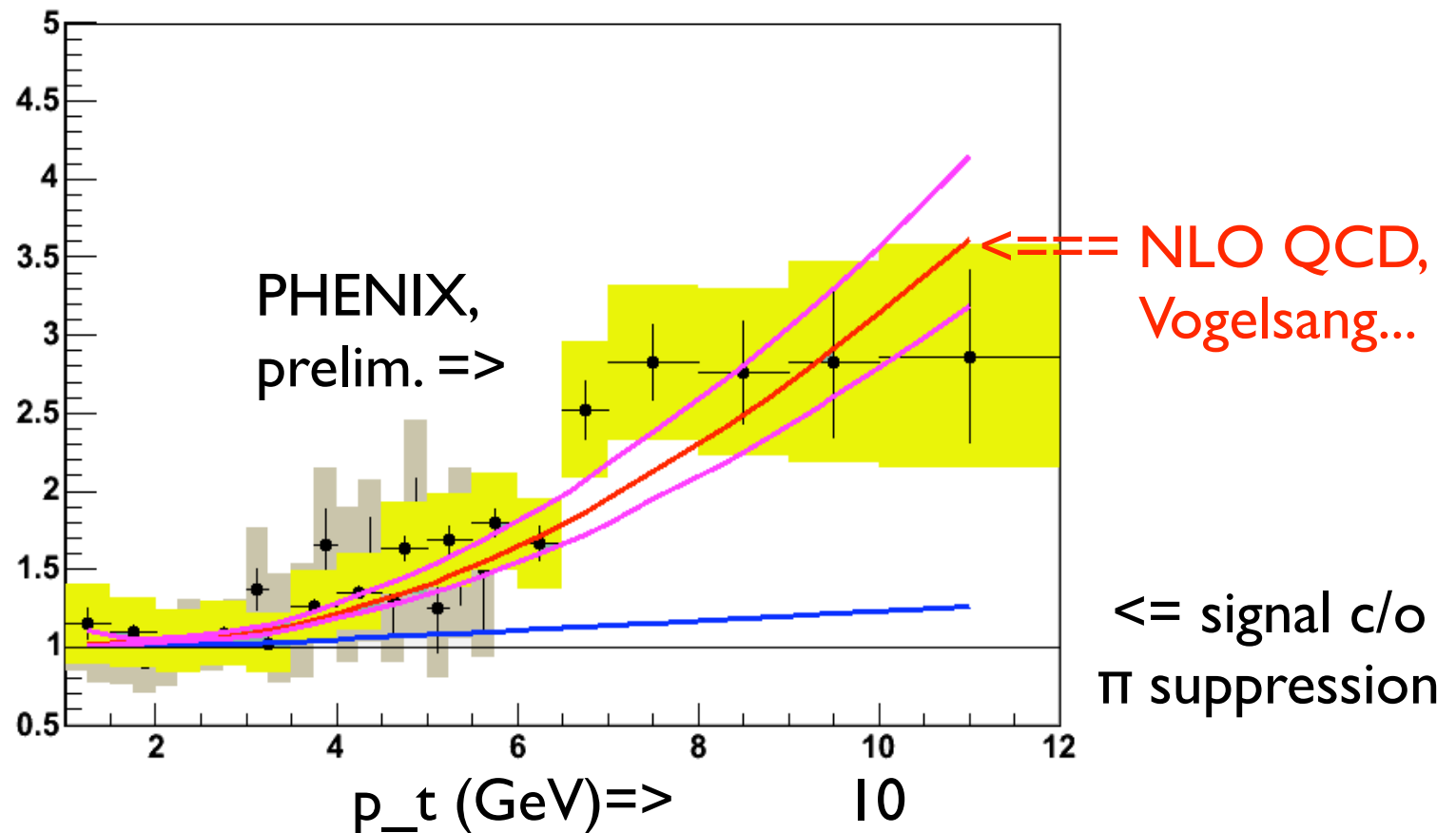
In direction *along* jet, *more* particles at low $p_{t,jet}$ in central AA than pp.

=> “stuff” in central AA shifts backward jet to *low* momentum,
forward jet *drags* “stuff” along with it!



Direct Photons Measured

Direct photons: easily escape, so probe initial state. *Without* pion suppression, very hard to measure (true at SPS). *With* observed suppression of π^0 's, measurable. Reasonable agreement at $p_t \sim 10$ GeV with Next to Leading Order QCD calculation, = pp times # binary collisions.



Has RHIC found (tamed) the “Unicorn” = QGP?

New final state effects:

R_AA

Suppression of backward jets

Also: new initial state effects,

Color Glass in forward dA

Exp'y: for the unicorn of central AA,
the high p_t “tail” wags the
low p_t “body”

HBT? Space-time evolution of the body?
Precise measure of thermal equilibration?
 p_t fluctuations at low p_t

Perhaps: it is a different beast....
But its still a *NEW* beast!





"A possible eureka."